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**Key Points:**

- Hydration reduces  $K_0$  of wadsleyite significantly
- Hydration does not influence  $dK/dP$  of wadsleyite
- Hydration reduces bulk sound velocity at transition zone conditions

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## Comparative compressibility of hydrous wadsleyite and ringwoodite: Effect of $H_2O$ and implications for detecting water in the transition zone

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**Abstract** Review of recent mineral physics literature shows consistent trends for the influence of Fe and  $H_2O$  on the bulk modulus ( $K_0$ ) of wadsleyite and ringwoodite, the major phases of Earth's mantle transition zone (410–660 km). However, there is little consensus on the first pressure derivative,  $K_0' = (dK/dP)_{P=0}$ , which ranges from about 4 to >5 across experimental studies and compositions. Here we demonstrate the importance of  $K_0'$  in evaluating the bulk sound velocity of the transition zone in terms of water content and provide new constraints on the effect of  $H_2O$  on  $K_0'$  for wadsleyite and ringwoodite by conducting a comparative compressibility study. In the experiment, multiple crystals of hydrous  $FO_{90}$  wadsleyite containing 2.0 and 0.25 wt %  $H_2O$  were loaded into the same diamond anvil cell, along with hydrous ringwoodite containing 1.4 wt %  $H_2O$ . By measuring their pressure-volume evolution simultaneously up to 32 GPa, we constrain the difference in  $K_0'$  independent of the pressure scale, finding that  $H_2O$  has no effect on  $K_0'$ , whereas the effect of  $H_2O$  on  $K_0$  is significant. The fitted  $K_0'$  values of hydrous wadsleyite (0.25 and 2.0 wt %  $H_2O$ ) and hydrous ringwoodite (1.4 wt %  $H_2O$ ) examined in this study were found to be identical within uncertainty, with  $K_0' \sim 3.7(2)$ . New secondary-ion mass spectrometry measurements of the  $H_2O$  content of these and previously investigated wadsleyite samples shows the bulk modulus of wadsleyite is reduced by 7.0(5) GPa/wt %  $H_2O$ , independent of Fe content for upper mantle compositions. Because  $K_0'$  is unaffected by  $H_2O$ , the reduction of bulk sound velocity in very hydrous regions of transition zone is expected to be on the order of 1.6%, which is potentially detectable in high-resolution, regional seismology studies.

### 1. Introduction

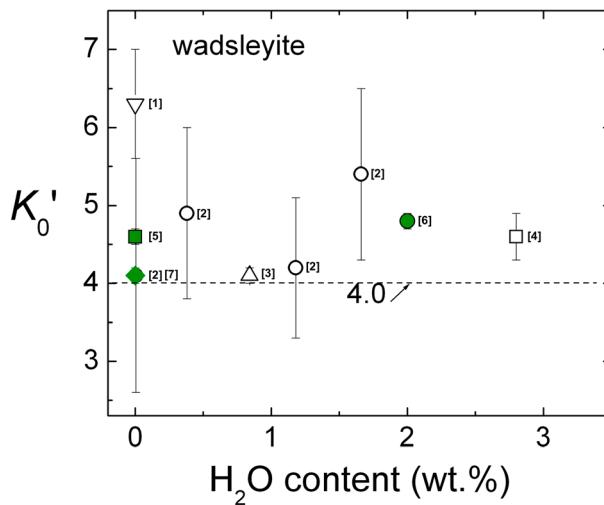
Interpreting the internal structure, composition, and mineralogy of the Earth's mantle requires accurate thermoelastic properties of the major mineral phases and critically how these properties vary with pressure, temperature, and composition [Birch, 1952; Duffy and Anderson, 1989; Li and Liebermann, 2007; Stixrude and Lithgow-Bertelloni, 2005]. Characteristics of the seismic discontinuities at 410 km and 660 km bounding the mantle transition zone, including the depth to transition, transition interval, and impedance contrast, are affected by the chemical composition of major mantle minerals [Bina and Helffrich, 1994; Frost, 2003; Irfune and Isshiki, 1998; Schmerr and Garnero, 2007; Sinogeikin et al., 2003]. Wadsleyite ( $\beta\text{-Mg}_2\text{SiO}_4$ ) and ringwoodite ( $\gamma\text{-Mg}_2\text{SiO}_4$ ) constitute ~60% of the transition zone in the pyrolite model of mantle composition [Irfune and Ringwood, 1987; Ringwood and Major, 1967]. In  $(\text{Mg},\text{Fe})_2\text{SiO}_4$ , upper mantle olivine typically contains  $X_{\text{Fe}} = \text{Fe}/(\text{Fe} + \text{Mg})$  of 0.07–0.15, making it essential to study Fe-bearing phases with well-characterized  $\text{Fe}^{3+}/\Sigma\text{Fe}$  ratios. The recent discovery of 1.5 wt %  $H_2O$  in a natural, Fe-bearing ringwoodite inclusion in diamond [Pearson et al., 2014] suggests that some regions of the transition zone may be near saturated in  $H_2O$ . Determining the combined influence of Fe and hydration on the elastic properties of wadsleyite and ringwoodite will therefore further constrain the chemical composition and water budget in the Earth's interior by comparison of experimental equations of state with high-resolution, regional seismic velocity profiles of the upper mantle.

The effect of Fe on the elastic properties of anhydrous wadsleyite can be evaluated by comparing previous studies on Mg- and Fe-bearing wadsleyite using static compression, ultrasonic, or optical-acoustic methods [e.g., Hazen *et al.*, 2000a, 2000b; Li and Gwanmesia, 1996; Li and Liebermann, 2000; Liu *et al.*, 2009; Sawamoto *et al.*, 1984; Sinogeikin *et al.*, 1998; Wang *et al.*, 2014; Zha *et al.*, 1997]. Fe increases the density of wadsleyite by about 3.5% for  $X_{\text{Fe}} = 0.1$  compared with pure Mg-wadsleyite [Li and Liebermann, 2007] but has a negligible effect on the bulk modulus [Wang *et al.*, 2014]. Although there is good agreement between studies on the bulk modulus of anhydrous Fe-bearing wadsleyite, the reported pressure derivative,  $K_0' = (dK/dP)_{P=0}$ , ranges from 4.1(1) [Wang *et al.*, 2014] to 4.6(1) [Li and Liebermann, 2000] or 4.8 [Gwanmesia *et al.*, 1990]. The range of experimental values for  $K_0'$  of wadsleyite is around 5 times the reported experimental uncertainties.

$\text{H}_2\text{O}$  is another important geochemical variable in mantle composition. Wadsleyite incorporates water as hydroxyl ( $\text{OH}^-$ ) defects into its structure due to the underbonded, nonsilicate oxygen ( $\text{O}1$ ) [Smyth, 1987, 1994], with up to  $\sim 3$  wt %  $\text{H}_2\text{O}$  reported in synthetic samples [Bofan-Casanova *et al.*, 2000; Demouchy *et al.*, 2005; Inoue *et al.*, 1995; Kohlstedt *et al.*, 1996]. Due to coupled substitutions with Mg-Fe vacancies on the M3 site [Jacobsen *et al.*, 2005; Deon *et al.*, 2010; Griffin *et al.*, 2013], the incorporation of hydroxyl in the wadsleyite structure has significant effect on its elastic properties. Previous studies on Mg-wadsleyite show that its bulk modulus decreases by  $\sim 12$  GPa per wt %  $\text{H}_2\text{O}$  added [Holl *et al.*, 2008; Mao *et al.*, 2008, 2011; Ye *et al.*, 2010], which is also in accordance with theoretical studies [Tsuchiya and Tsuchiya, 2009; Liu *et al.*, 2012]. Whereas previous studies of the bulk modulus have investigated either hydrous Mg-wadsleyite or anhydrous Fe-bearing wadsleyite, the combined influence of hydration and iron on the elasticity of wadsleyite at high pressure has been reported in only one experimental study thus far [Mao *et al.*, 2011], who found the effects of Fe and  $\text{H}_2\text{O}$  on the bulk and shear moduli are linearly additive and reported  $K_0' = 4.8(1)$  for  $X_{\text{Fe}} = 0.10$  hydrous wadsleyite containing about 2 wt %  $\text{H}_2\text{O}$ . The hydrous Fe-bearing wadsleyite sample from the study of Mao *et al.* [2011] is also used in the current study for comparison with low- $\text{H}_2\text{O}$ , Fe-bearing wadsleyite. A recent theoretical study [Tian *et al.*, 2012] on the elasticity of hydrous Fe-bearing wadsleyite predicted that the incorporation of 1.6 wt %  $\text{H}_2\text{O}$  into Fe-bearing wadsleyite ( $X_{\text{Fe}} = 0.125$ ) elevates  $K_0'$  from 3.92 to 4.12.

Whereas the effect of hydration on the bulk modulus of hydrous wadsleyite and ringwoodite is generally well constrained at room pressure and room temperature [e.g., Jacobsen, 2006], the influence of  $\text{H}_2\text{O}$  on  $K_0'$  is poorly constrained when considering the range of reported values cited in the literature. The pressure derivative of  $K$  is critical for estimating the effect of  $\text{H}_2\text{O}$  on the density and sound velocities of wadsleyite at transition zone pressures. The influence of hydration on the pressure derivative of the bulk modulus for wadsleyite and ringwoodite may not be well constrained for various reasons, including limited data resolution at high pressure, the use of different pressure calibrations, nonhydrostatic stresses on the sample, different data analysis methods, or the narrow pressure ranges over which data have been measured. Reported values of the bulk modulus of anhydrous Mg-wadsleyite ( $\text{Fo}_{100}$ ) are in good agreement with  $K_0 \sim 170$  GPa; however, values of  $K_0'$  range from 4.2 to 6.3 [e.g., Holl *et al.*, 2008; Hazen *et al.*, 2000a; Li and Gwanmesia, 1996; Zha *et al.*, 1997]. Figure 1 summarizes the  $K_0'$  values of wadsleyite from published experimental studies. The values of  $K_0'$  range from 4 to almost 7 and do not show a systematic dependence on  $\text{H}_2\text{O}$  content. Extrapolating the equation of state of anhydrous Mg-wadsleyite to 20 GPa by using  $K_0'$  of 5 instead of 4 would result a 0.5% difference in density, which is significant in seismic observations. Moreover,  $K_0$  and  $K_0'$  derived from least squares fitting  $P$ - $V$  data sets have intrinsic dependence between fitted variables, i.e., anticorrelation between  $K_0$  and  $K_0'$  [Bass *et al.*, 1981; Bell *et al.*, 1987; Bina, 1995]. This anticorrelation, often ignored when comparing results across experiments, must be considered carefully when analyzing the effects of chemical composition on elastic properties.

In this study, we focus on accurately determining the effect of hydration on the difference in  $K_0'$  between nearly anhydrous and hydrous Fe-bearing wadsleyite. By simultaneously loading wadsleyite crystals of similar Fe content but significantly different  $\text{H}_2\text{O}$  contents into one diamond anvil cell experiment, we measured their relative volumes at each pressure and a direct comparison of  $K_0'$  is made without introducing bias from the pressure scales. Furthermore, the data sets are equivalent in that every volume measurement for each composition shares a common pressure point. In the comparative compressibility experiments, three Fe-bearing ( $X_{\text{Fe}} = 0.10$ ) hydrous wadsleyite crystals containing 2.0 wt %  $\text{H}_2\text{O}$  were loaded together with three crystals having the same Fe content but only 0.25 wt %  $\text{H}_2\text{O}$ , an order of magnitude less water. In addition, we included a crystal of hydrous ringwoodite that was synthesized coexisting with the hydrous wadsleyite in order to constrain



**Figure 1.** Pressure derivative of the bulk modulus,  $K_0' = (dK/dP)_{P=0}$ , versus H<sub>2</sub>O content for Mg-wadsleyite (open symbols) and Fe-bearing wadsleyite (filled symbols) compiled from the experimental literature. References are [1] Hazen *et al.* [2000a], [2] Holl *et al.* [2008], [3] Mao *et al.* [2008], [4] Yusa and Inoue [1997], [5] Li and Liebermann [2000], [6] Mao *et al.* [2011], and [7] Wang *et al.* [2014].

the difference in  $K_0'$  between hydrous ringwoodite and hydrous wadsleyite. A highly focused synchrotron X-ray beam of  $\sim 10 \mu\text{m}$  diameter allowed the collection of individual diffraction patterns from the seven different crystals within the same diamond anvil cell. In addition, two internal pressure standards (MgO and ruby fluorescence) were employed and small pressure steps of  $< 1 \text{ GPa}$  up to  $32 \text{ GPa}$  facilitated fitting and analysis of a third-order Birch-Murnaghan equation of state to the pressure-volume data. The results constrain the relative value of  $K_0'$  between dry and hydrous wadsleyite and also between the coexisting hydrous Fe-bearing wadsleyite and ringwoodite. These results are used to evaluate the effect of iron and hydrogen on the bulk modulus of wadsleyite and ringwoodite at pressures relevant to the mantle transition zone.

## 2. Experimental Methods

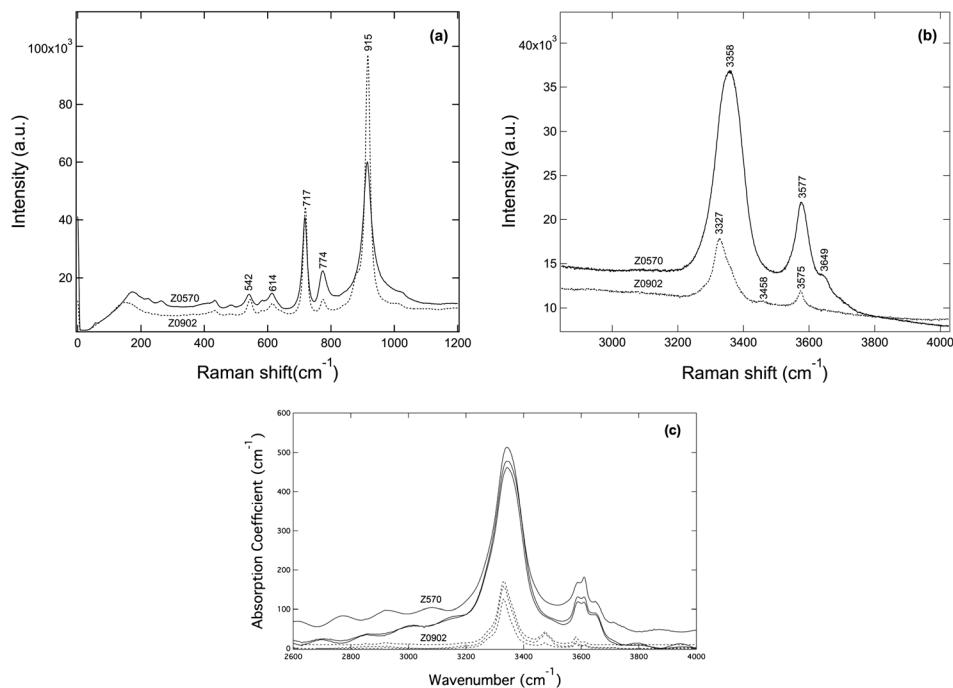
### 2.1. Sample Synthesis

Single crystals of Fe-bearing hydrous wadsleyite and ringwoodite were synthesized in the 5000 t, multianvil press at Bayerisches Geoinstitut (run Z0570). Crystals of hydrous wadsleyite used in the high-pressure Brillouin spectroscopy study of Mao *et al.* [2011] came from the same synthesis run. Starting materials consisted of San Carlos olivine plus  $\sim 10 \text{ wt \% H}_2\text{O}$  added as liquid and welded into a Pt capsule 3 mm in height and diameter. The starting material was compressed to 18 GPa and heated to 1400°C. Heating duration was 1 min at 1400°C and then decreased to 1100°C and held for 20 min before quenching. The anhydrous wadsleyite (run Z0920) was synthesized from San Carlos olivine at 11 GPa and 1100°C using the same technique as run Z0570, but without the addition of H<sub>2</sub>O.

### 2.2. Raman and FTIR Spectroscopy

Wadsleyite crystals from runs Z0570 and Z0920 were initially characterized by Raman spectroscopy. Unpolarized Raman spectra were obtained using a solid-state excitation laser with  $\lambda_0 = 458 \text{ nm}$  and recorded from 120 to  $4000 \text{ cm}^{-1}$ . In the low-wave-number region, the hydrous (Z0570) and nominally anhydrous (Z0920) wadsleyite exhibit similar Raman spectra, both in the number and frequency of major vibrational modes (Figure 2a). At high wave number (Figure 2b), the O-H stretching modes are observed in both samples, indicating that some H<sub>2</sub>O is incorporated into the nominally anhydrous sample Z0920 even though no water was intentionally added to the experiment. Although one cannot quantify the differences in intensity of the O-H stretching modes from unpolarized Raman spectra, the data shown in Figure 2b were collected with identical parameters. The intensity of the main O-H stretching band at  $3358 \text{ cm}^{-1}$  for the hydrous sample Z0570 is roughly 3 times higher than for the nominally anhydrous sample Z0920. The band is also shifted to lower wave numbers by about  $30 \text{ cm}^{-1}$  for the less hydrous sample, Z0920. A similar difference of intensity is observed in the second strongest O-H stretching band at  $3576 \text{ cm}^{-1}$ , but there is no difference in shift between the two samples (Figure 2b). Sample Z0570 displays a weak band at  $3649 \text{ cm}^{-1}$ , which is not observed in sample Z0920.

In order to further evaluate the difference in H<sub>2</sub>O content between wadsleyite samples Z0570 and Z0920, unpolarized Fourier transform infrared (FTIR) spectra were obtained on three different randomly oriented crystals from each run (Figure 2c). Crystals were double polished with a final thickness of  $18 \pm 2 \mu\text{m}$  for all



**Figure 2.** Unpolarized Raman and FTIR spectra of Fe-bearing wadsleyite Z0902 (dashed lines, 0.25 wt % H<sub>2</sub>O) and Z0570 (solid lines, 2.0 wt % H<sub>2</sub>O). (a) Raman spectra in the lattice mode region. (b) Raman spectra in the O-H stretching region. (c) Unpolarized FTIR spectra from three different crystals (random orientation) of each composition. All three crystals of Z0570 were 18 ± 2 μm thick (solid lines), giving a maximum absorbance of ~0.9. The crystals of Z0902 were 21 ± 2, 21 ± 2, and 46 ± 2 μm thick, going from the top to bottom spectrum (dashed lines).

samples from Z570, and for samples from Z0902 the three crystals were 21 ± 2, 21 ± 2, and 46 ± 2 μm thick. The maximum absorbance (for Z570) was 0.9. IR spectra were collected from 1850 to 4000 cm<sup>-1</sup> with a globar light source, Bruker IFS 66v FTIR spectrometer, Hyperion microscope, and InSb detector. Following baseline correction, integrated absorbance and area-weighted averages of the peak positions were used to estimate the H<sub>2</sub>O content of each sample using the wadsleyite-specific calibration of Deon *et al.* [2010]. By this method, we obtain 1.4 wt % H<sub>2</sub>O for sample Z0570 and 0.25 wt % H<sub>2</sub>O for sample Z0902.

### 2.3. Secondary-Ion Mass Spectrometry

The H<sub>2</sub>O content of wadsleyite and ringwoodite from run Z0570 was also measured by secondary-ion mass spectrometry (SIMS). We present an interlaboratory comparison made by measuring the H<sub>2</sub>O content of wadsleyite Z0570 crystals on two different instruments: first on the Cameca IMS-3F ion mass microanalyzer at Hokkaido University and subsequently on the Cameca NanoSIMS 50L at the Carnegie Institution of Washington, Department of Terrestrial Magnetism (DTM). Table 1 gives experimental parameters and results of the interlaboratory SIMS measurements. Briefly, on the IMS-3F instrument at Hokkaido University, five different crystals measuring 200 μm across or smaller were measured with at least two points per crystal. In the Hokkaido measurements, a natural amphibole crystal containing 1.66 wt % H<sub>2</sub>O was used as a calibration standard, as described by Miyagi *et al.* [1998]. At Carnegie Institution on the NanoSIMS, three different silicate glasses containing 0.17, 1.00, and 1.62 wt % H<sub>2</sub>O were used to develop a calibration against the measured mass ratios of <sup>17</sup>OH/<sup>28</sup>Si [Hauri *et al.*, 2006; Saal *et al.*, 2008]. Further experimental details are provided in Table 1. The average H<sub>2</sub>O content measured from 5 different crystals and 19 different measurement points at Hokkaido is 2.07(± 0.24) wt % H<sub>2</sub>O. The Hokkaido SIMS measurements are in agreement within mutual experimental uncertainty with NanoSIMS results from DTM-Carnegie, where the average of 14 measurements points across one crystal was 1.93(± 0.22) wt % H<sub>2</sub>O. We take the average of the interlaboratory results, 2.0(2) wt % H<sub>2</sub>O, to be the best estimate of the water content of hydrous wadsleyite Z0570. The SIMS H<sub>2</sub>O content of wadsleyite Z0570 is about 1.4 times (roughly 2–3σ) higher than suggested by FTIR using the Deon *et al.* [2010]

**Table 1.** Interlaboratory Comparison of Secondary-Ion Mass Spectrometry (SIMS) Parameters and Results for the H<sub>2</sub>O Content Wadsleyite, Sample Z0570

Parameter	Experiment 1	Experiment 2
Facility	Hokkaido University	Carnegie DTM
Instrument	Cameca IMS-3F	Cameca NanoSIMS 50L
Beam	<sup>16</sup> O <sup>-</sup>	Cs <sup>+</sup>
Voltage/current	12.5 kV/10 nA	8 kV/1 nA
Beam or raster size	30 μm diameter (beam)	7 × 7 μm <sup>2</sup> (raster)
Ion collection area	10 × 10 μm <sup>2</sup>	2.6 × 2.6 μm <sup>2</sup>
Measured ratio	<sup>1</sup> H <sup>+</sup> / <sup>30</sup> Si <sup>+</sup>	<sup>17</sup> OH/ <sup>28</sup> Si
Standards	Amphibole (1.66 wt % H <sub>2</sub> O)	Silicate glasses
Detection limit sample	San Carlos olivine	Synthetic forsterite
Measured value	0.0049 wt % H <sub>2</sub> O	0.0025 wt % H <sub>2</sub> O
H <sub>2</sub> O content (wt % H <sub>2</sub> O)	Crystal 1 2.11 1.90 1.78 2.65 1.81 2.41 Crystal 2 2.07 1.97 1.85 Crystal 3 2.31 1.83 Crystal 4 2.15 2.03 2.30 2.22 Crystal 5 1.96 1.83 1.87 2.27	Crystal 6 1.74 1.75 1.86 1.88 1.88 2.04 2.35 2.32 1.64 1.56 2.01 2.04 1.90 1.99
Average wt % H <sub>2</sub> O	2.07(24) wt % H <sub>2</sub> O	1.93(22) wt % H <sub>2</sub> O

calibration. Given the importance of hydrous wadsleyite to modeling the mineralogy and geochemistry of Earth's transition zone, future work on calibrating both spectroscopic and SIMS H<sub>2</sub>O contents for wadsleyite against absolute methods is warranted [e.g., Thomas *et al.*, 2015].

For the hydrous ringwoodite synthesized in run Z0570 and coexisting with hydrous wadsleyite, NanoSIMS measurements at DTM-Carnegie give an average from nine measurements across one crystal of 1.38 ( $\pm 0.30$ ) wt % H<sub>2</sub>O. The ratio of H<sub>2</sub>O in wadsleyite to coexisting ringwoodite in experiment Z0570 is thus  $\sim 1.5$ , which is comparable to the hydrogen partitioning value of  $\sim 2$  between wadsleyite and ringwoodite reported by Inoue *et al.* [2010]; in both cases wadsleyite contains more H<sub>2</sub>O than the coexisting ringwoodite.

#### 2.4. Major Element Chemistry: EMPA and Mössbauer Spectroscopy

The major element composition of wadsleyite and ringwoodite samples was determined by electron microprobe analysis (EMPA). The resulting Mg/(Mg + Fe<sup>2+</sup> + Fe<sup>3+</sup>) ratios were 0.90 and 0.89 for wadsleyite samples in run Z0902 and run Z0570, respectively. Samples from run Z0902 and run Z0570 both contain one Si per formula unit. The oxidation state of Fe in hydrous wadsleyite (run Z0570) was characterized by Mössbauer spectroscopy. A piece of 25 μm thick Ta foil (absorbing 99% of 14.4 keV gamma rays) drilled with either a 300 μm hole (sample 1) or 400 μm hole (sample 2) was centered over each crystal. The average thickness of the

crystals was estimated to be roughly 200  $\mu\text{m}$  (sample 1) and 250  $\mu\text{m}$  (sample 2), respectively, which combined with the composition of  $\text{Fo}_{90}$  (i.e.,  $X_{\text{Fe}}=0.10$ ) implies an effective Mössbauer thickness of 6.9 and 8.6 mg Fe/cm<sup>2</sup> for samples 1 and 2, respectively. The Mössbauer spectra were recorded at room temperature (293 K) in transmission mode on a constant acceleration Mössbauer spectrometer with a nominal 370 MBq <sup>57</sup>Co high specific activity source in a 12  $\mu\text{m}$  thick Rh matrix. The velocity scale was calibrated relative to a 25  $\mu\text{m}$  thick  $\alpha$ -Fe foil using the positions certified for (former) National Bureau of Standards standard reference material no. 1541; line widths of 0.36 mm/s for the outer lines of  $\alpha$ -Fe were obtained at room temperature. The spectra were collected over 2 days (sample 1) and 1 day (sample 2), respectively, and were fitted using the commercially available fitting program NORMOS written by R.A. Brand (distributed by Wissenschaftliche Elektronik GmbH, Germany).

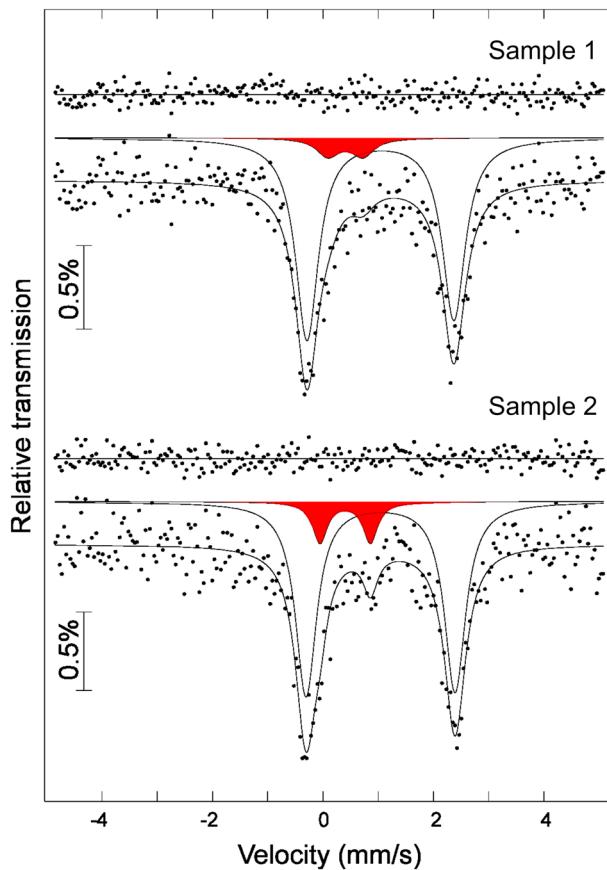
The Mössbauer spectra of hydrous Fe-bearing wadsleyite from run Z0570 are dominated by an asymmetric quadrupole doublet with less intense absorption between the main components (Figure 3). The dominant absorption was fitted to a single  $\text{Fe}^{2+}$  doublet, and the remaining absorption was modeled using a doublet corresponding to  $\text{Fe}^{3+}$ . Conventional constraints were applied to the components of each  $\text{Fe}^{3+}$  doublet (i.e., equal widths and areas), while only the widths of the  $\text{Fe}^{2+}$  components were constrained to be equal. The area ratios of the  $\text{Fe}^{2+}$  doublet components were allowed to vary, due to the effects of preferred orientation arising from the single-crystal nature of the sample. In addition, the center shift of the  $\text{Fe}^{3+}$  doublet from sample 1 was constrained to the value obtained from the spectrum for sample 2. All spectral components were assumed to have Lorentzian line shape. Estimated standard deviations were derived from the statistics of the fitting process as well as uncertainties in the fitting model itself. The hyperfine parameters derived from the fits are listed in Table 2.

The center shifts and quadrupole splitting values for the  $\text{Fe}^{2+}$  and  $\text{Fe}^{3+}$  doublets are consistent within experimental uncertainty with values previously observed for hydrous wadsleyite [Bofan-Casanova *et al.*, 2012]. The values of  $\text{Fe}^{3+}/\Sigma\text{Fe}$  were calculated from the relative areas of the  $\text{Fe}^{3+}$  absorption. The resulting average of  $\text{Fe}^{3+}/\Sigma\text{Fe}$  in sample Z0570 is about 10%. Although the ratio  $\text{Fe}^{3+}/\Sigma\text{Fe}$  of the sample in run Z0902 was not characterized, it is expected to be very low because of the low water content in this sample [McCammon *et al.*, 2004]. After the chemical compositions were known, the chemical formulae of the samples in run Z0902 and run Z0570 were determined to be  $(\text{Mg}_{1.782}^{2+}\text{Fe}_{0.198}^{2+})\text{H}_{0.041}\text{SiO}_4$  and  $(\text{Mg}_{1.638}^{2+}\text{Fe}_{0.173}^{2+}\text{Fe}_{0.019}^{3+})\text{H}_{0.320}\text{SiO}_4$  based on charge balance, water content, Mg/(Mg +  $\text{Fe}^{2+}$  +  $\text{Fe}^{3+}$ ) ratio, and  $\text{Fe}^{3+}/\Sigma\text{Fe}$  ratio. The same chemical analyses were applied to the Fe-bearing ringwoodite (Z0570) sample used in this experiment, yielding the chemical formula  $(\text{Mg}_{1.577}^{2+}\text{Fe}_{0.291}^{2+}\text{Fe}_{0.032}^{3+})\text{H}_{0.227}\text{Si}_{0.985}\text{O}_4$ .

## 2.5. High-Pressure Single-Crystal X-Ray Diffraction

Before carrying out the high-pressure experiment, initial lattice parameters of samples under ambient conditions were measured using a four-circle X-ray diffractometer at the University of Colorado Boulder. Single-crystal lattice parameters were refined using the program SINGLE [Angel and Finger, 2011]. The low- $\text{H}_2\text{O}$  wadsleyite sample (run Z0902) is orthorhombic, whereas hydrous wadsleyite (Z0570) is slightly monoclinic with  $\beta=90.078(2)^\circ$ . Room pressure lattice parameters of the hydrous wadsleyite were also refined with orthorhombic symmetry for comparison with subsequent high-pressure X-ray diffraction (XRD) data collected at the synchrotron. Lattice parameters, density, and molar volume of the samples used in this study are reported in Table 3.

High-pressure, single-crystal X-ray diffraction was carried out at beamline 16-ID-B, HPCAT, Advance Photon Source, Argonne National Laboratory. The wavelength of the monochromatic synchrotron radiation X-ray was  $\lambda=0.407376 \text{ \AA}$ , calibrated using a  $\text{CeO}_2$  standard. Diffraction images were recorded by a MAR CCD detector, positioned 21.32 cm from the sample. A symmetric-piston diamond cell, fitted with diamond anvils having 300  $\mu\text{m}$  culets, was used. A rhenium gasket was preindented to about 30  $\mu\text{m}$  initial thickness, and a hole approximately 200  $\mu\text{m}$  in diameter was eroded using an electrostatic discharge machine. The diamond anvil cell was also fitted with one cubic boron nitride seat, placed downstream of the synchrotron X-ray source to allow wide-angle access to diffracted X-rays, while a slit tungsten carbide seat was used in the upstream side. This setup allowed oscillation of the cell about the vertical axis by  $\pm 14^\circ$ . Three single-crystal anhydrous wadsleyites (run Z0902) as well as three single-crystal hydrous wadsleyites (run Z0570) and one crystal of hydrous ringwoodite (run Z0570) were loaded together into the same diamond anvil cell, as shown in Figure 4.



**Figure 3.** Mössbauer spectra of two different hydrous wadsleyite crystals (labeled sample 1 and sample 2) from run Z0570. The absorption assigned to  $\text{Fe}^{3+}$  is indicated in red and constitutes an average of about 10% of the total Fe in the sample.

After increasing pressure at each pressure step, the cell was allowed to settle for 10–15 min before performing XRD to ensure that pressure was stable during measurements. Three annealed ruby spheres were loaded along with samples into the pressure chamber as pressure markers. The ruby sphere at the center of the cell served as the primary marker, while the other two at the edges were used for checking pressure gradients. The ruby pressure was calibrated from the measured shift ( $\lambda/\lambda_0$ ) of the R1 ruby fluorescence line using  $P$  (GPa) =  $A/B[(1 + (\Delta\lambda/\lambda_0))^B - 1]$ , where  $A = 1904$  and  $B = 10.32$  [Jacobsen et al., 2008]. The reported ruby pressure in this paper is the average of pressures measured before and after collecting XRD images, which did not differ by more than 0.2 GPa. A single-crystal of MgO was also loaded into the pressure chamber as a secondary pressure marker. MgO pressures were calculated using the third-order Birch-Murnaghan (BM) equation of state (EOS) with initial volume  $V_0 = 74.698$  ( $\text{\AA}^3$ ),  $K_{T0} = 160.32$  GPa, and  $K_0' = 4.08$  [Kennett and Jackson, 2009]. Uncertainties of MgO pressures were estimated following the procedures described in Angel et al. [1997].

**Table 2.** Results From Mössbauer Spectroscopy<sup>a</sup>, Wadsleyite Sample Z0570

	CS (mm/s) (Relative to $\alpha$ -Fe)	QS (mm/s)	FWHM (mm/s)	Area
Crystal 1				
$\text{Fe}^{2+}$	1.04(1)	2.66(5)	0.50(5)	0.92(5)
$\text{Fe}^{3+}$	0.4 <sup>b</sup>	0.63(16)	0.49(27)	0.08(5)
Crystal 2				
$\text{Fe}^{2+}$	1.04(1)	2.70(5)	0.46(5)	0.87(6)
$\text{Fe}^{3+}$	0.40(5)	0.91(10)	0.32(12)	0.13(6)

<sup>a</sup>CS = center shift; QS = quadrupole Splitting; FWHM = full width at half maximum.

<sup>b</sup>Constrained to the value obtained from crystal 2.

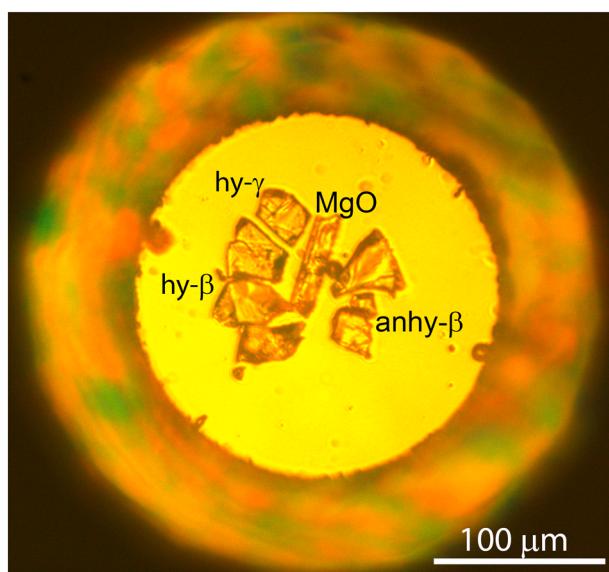
By loading three crystals, each of hydrous and anhydrous wadsleyite together in the same diamond anvil cell, we achieved better  $hkl$  coverage. The cell was loaded with neon as pressure medium using the COMPRES/GSECARS gas loading system at Advanced Photon Source [Rivers et al., 2008]. Two pressure standards were also included in the cell; a single-crystal of MgO and an annealed ruby sphere, for the purpose of comparing the results from two pressure scales on the fitted  $K_0'$ . The pressure inside the cell was about 1.09 GPa after gas loading, and the gasket hole shrank to 75% of its original size. Neon was chosen in this experiment for its ability to retain a quasi-hydrostatic state up to 80 GPa [Dewaele et al., 2008], which enables it to maintain similar pressure environments surrounding every crystal inside the cell.

Because  $K_0'$  is the second derivative of a  $P$ - $V$  curve, a precise determination of pressures and fine pressure steps are critical in constraining  $K_0'$  in a high-pressure compression experiment. In order to better constrain the equation of state, a gas membrane cell as well as an online ruby system was set up to control small pressure increments of 0.5–0.8 GPa per step up to 32 GPa.

**Table 3.** Composition, Lattice, and Density Parameters of Samples Used in This Study

Run #	Z0902- $\beta$ Phase	Z0570- $\beta$ Phase	Z0570- $\gamma$ Phase
Mg/ $\Sigma$ (Mg + Fe <sup>2+</sup> + Fe <sup>3+</sup> )	0.90	0.89	0.83
Fe <sup>3+</sup> / $\Sigma$ Fe (%)	very low	~10%	~10%
H <sub>2</sub> O content (wt %)	~0.25	2.0	1.38
Lattice parameters	$\beta_0 = 90^\circ$ $a_0 = 5.7045(4)$ Å $b_0 = 11.4765(4)$ Å $c_0 = 8.2707(4)$ Å $V_0 = 541.46(5)$ Å <sup>3</sup>	$\beta_0 = 90^\circ$ $a_0 = 5.6918(10)$ Å $b_0 = 11.5276(10)$ Å $c_0 = 8.2641(8)$ Å $V_0 = 542.23(12)$ Å <sup>3</sup> $\beta_0 = 90.078(2)^\circ$ $a_0 = 5.6929(2)$ Å $b_0 = 11.5275(2)$ Å $c_0 = 8.2648(1)$ Å $V_0 = 542.37(2)$ Å <sup>3</sup>	$a_0 = 8.1011(4)$ Å $V_0 = 531.65(8)$ Å <sup>3</sup>
Atomic mass, $M$ (g/mol)	146.48	143.45	147.82
Density, $\rho_0$ (g/cm <sup>3</sup> )	3.595	3.515	3.695
Molar volume, MV (mol/cm <sup>3</sup> )	40.75	40.81	40.00

To obtain the orientation matrix at the initial pressure of 1.09 GPa, an  $\omega$  step scan was performed on each crystal in the range from  $-14$  to  $14^\circ$ , with step size  $1^\circ$ , and exposure time 2 s per image. These resulting 29 images were used to calculate the unit cell parameters and determine the peak positions. The peak positions for all crystals were assigned using the software GSE\_ADA [Dera, 2007a, 2007b]. A wide  $\omega$  scan was performed thereafter at subsequent pressure steps in the range from  $-14$  to  $14^\circ$  with exposure time 28 s on each crystal. For each composition, the unit cell parameters were refined by combining peaks collected from all three crystals in unit cell [Holland and Redfern, 1997]. The number of available peaks decreased with increasing pressure. In order to maintain consistency, only the peaks appearing in all pressure steps were included for high-pressure lattice parameter refinement. Thus, 52 diffracted peaks for sample Z0902 and 69 diffracted peaks for sample Z0570 were included in the unit cell analyses. Although there were only six diffracted peaks available for the hydrous ringwoodite sample, they were sufficient for the purpose of determining the cubic lattice parameters.

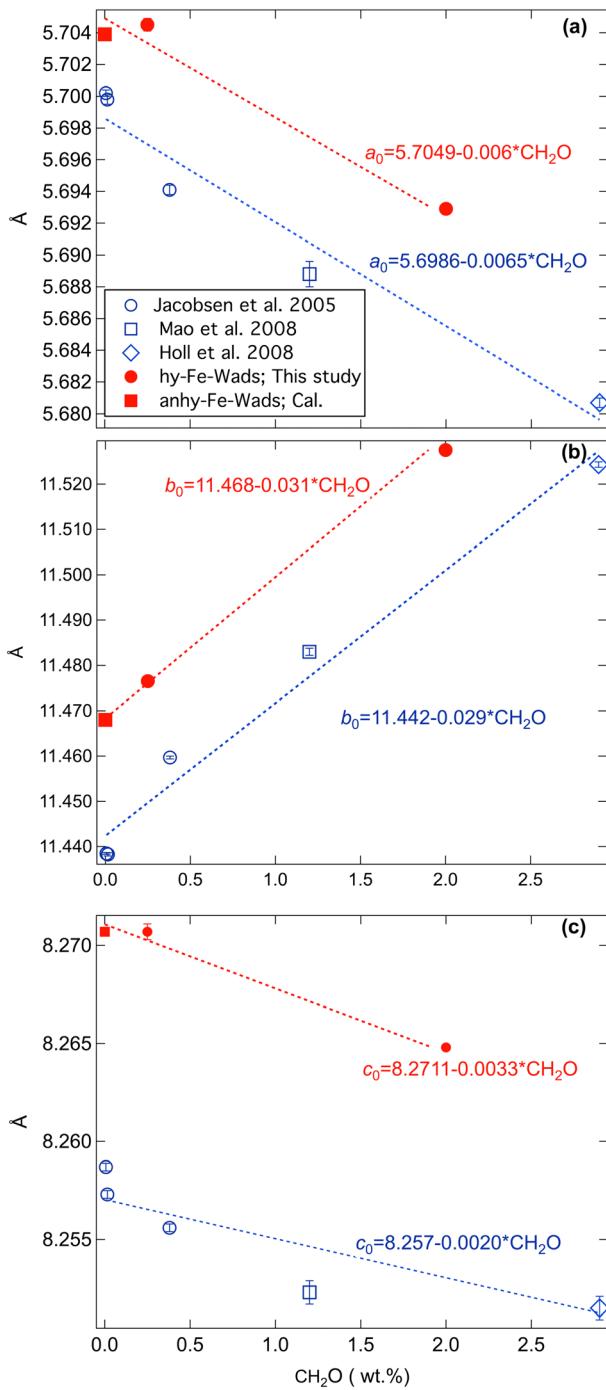


**Figure 4.** Diamond anvil cell setup for the comparative compressibility study, at room pressure before loading with neon. There are three single crystals of wadsleyite from run Z0902 (labeled anhy- $\beta$ ), three single crystals of wadsleyite from run Z0570 (labeled hy- $\beta$ ), and a crystal of coexisting ringwoodite from run Z0570 (labeled hy- $\gamma$ ). A single-crystal MgO and three ruby spheres were also included for pressure calibration.

### 3. Results and Discussion

#### 3.1. Influence of Fe and H<sub>2</sub>O on Lattice Parameters and Density of Wadsleyite

Whereas Fe increases all three lattice parameters of wadsleyite, H<sub>2</sub>O decreases  $a$  and  $c$  but increases the  $b$  axis, as illustrated in Figure 5. The Fe-bearing wadsleyite sample Z0902 (0.25 wt % H<sub>2</sub>O) lattice parameters are  $a_0 = 5.7045(4)$ ,  $b_0 = 11.2707(4)$ , and  $c_0 = 8.2707(4)$  Å, compared with the very hydrous (2.0 wt % H<sub>2</sub>O) and slightly monoclinic sample Z0570 with  $a_0 = 5.6929(2)$ ,  $b_0 = 11.5275(2)$ ,  $c_0 = 8.2648(1)$  Å, and  $\beta = 90.078(2)^\circ$ . Refining the lattice parameters of sample Z0570 in the orthorhombic system gives  $a_0 = 5.6918(10)$ ,  $b_0 = 11.5276(10)$ , and  $c_0 = 8.2641(8)$  Å. The unit cell refinements in this study further confirm that hydration lowers the symmetry of wadsleyite from  $I\bar{m}ma$  to  $I2/m$  [Smyth et al., 1997; Holl et al., 2008; Jacobsen et al., 2005;



**Figure 5.** Variation of lattice parameters with  $\text{H}_2\text{O}$  content in Mg-wadsleyite (blue) and Fe-bearing wadsleyite (red). The dash lines represent linear regressions.

determined by the  $b/a$  axial ratio calibration of Jacobsen *et al.* [2005], which was tied largely to Paterson [1982]. The FTIR water contents of the current samples, Z0570 and Z0920, are reported from the mineral-specific calibration of Deon *et al.* [2010]. For sample Z0570, the maximum absorbance was around 0.9 and we found for a mean wave number of  $3380 \text{ cm}^{-1}$ , the Deon *et al.* [2010] calibration leads to a 1.2 times higher water content than Libowitzky and Rossman [1997] and 1.6 times higher value than Paterson [1982]. The  $\text{H}_2\text{O}$  contents measured previously by FTIR and the new NanoSIMS results are listed in Table 4 and plotted in Figure 6.

Kudoh *et al.*, 1996]. Violation of orthorhombic symmetry can be attributed to the ordering of divalent cation vacancies onto nonequivalent M3 sites, breaking the mirror symmetry element perpendicular to the  $b$  axis [Smyth *et al.*, 1997].

To better constrain the combined effects of Fe and  $\text{H}_2\text{O}$  on wadsleyite lattice parameters, we first quantified the effects of Fe and  $\text{H}_2\text{O}$  separately, based on this and previous studies using data listed in Table 4 and plotted in Figure 5. Finger *et al.* [1993] investigated the effects of Fe on the structure of wadsleyite and showed that Fe prefers the M3 and M1 sites over the M2 site. The effect of iron on lattice parameters estimated from the anhydrous Fe-bearing wadsleyite data [Finger *et al.*, 1993] shows that increasing  $X_{\text{Fe}}$  by 0.10 increases  $a_0$ ,  $b_0$ , and  $c_0$  by 0.011, 0.017, and 0.016 Å, respectively. Analyzing the effects of hydration on Mg- and Fe-bearing wadsleyite requires a systematic approach to determining the water content.

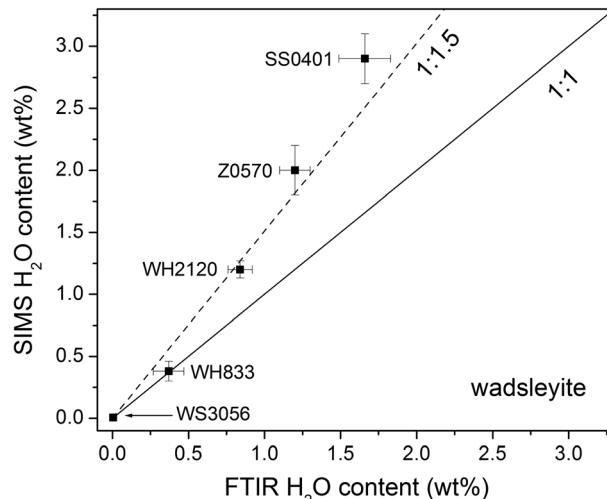
In order to eliminate issues associated with comparing samples characterized by different methods (FTIR, Raman, or SIMS), we redetermined the water content of several Mg-wadsleyite samples from the same synthesis runs as previous FTIR and elasticity studies using NanoSIMS. We followed the same procedure described above for sample Z0570. In particular, we analyzed the water content of WS3056 and WH833 from the study of Jacobsen *et al.* [2005], sample WH2120 from Mao *et al.* [2008], and sample SS0401 from the study of Holl *et al.* [2008]. Referring to Table 4, the  $\text{H}_2\text{O}$  content of WS3056 was previously determined by the calibration of Paterson [1982]. The  $\text{H}_2\text{O}$  content of WH833 was previously determined by the calibration of Libowitzky and Rossman [1997]. The  $\text{H}_2\text{O}$  contents of WH2120 and SS0401 were previously

**Table 4.** Compilation of Fe and H<sub>2</sub>O Content by Different Methods and Lattice Parameters of Hydrous Wadsleyite

Sample	$X_{\text{Fe}}$	H <sub>2</sub> O (wt %; FTIR)	H <sub>2</sub> O (wt %; SIMS)	$a$ (Å)	$b$ (Å)	$c$ (Å)	$V$ (Å <sup>3</sup> )	$\rho_0$ (kg/m <sup>3</sup> )	Reference
WS3056	0	0.005 <sup>a</sup>	0.006(3) <sup>f</sup>	5.7002(2)	11.4385(2)	8.2587(2)	538.46(2)	3472	Jacobsen et al. [2005]
WZ304	0	0.015 <sup>a</sup>		5.6998(2)	11.4383(3)	8.2573(2)	538.34(3)	3473	Jacobsen et al. [2005]
WH833	0	0.37(1) <sup>b</sup>	0.38(8) <sup>f</sup>	5.6941(2)	11.4597(3)	8.2556(2)	538.70(3)	3455	Jacobsen et al. [2005]
WH2120	0	0.84(8) <sup>c</sup>	1.20(7) <sup>f</sup>	5.6888(6)	11.4830(8)	8.2523(6)	539.08(8)	3418	Mao et al. [2008]
SS0401	0	1.66(17) <sup>c</sup>	2.9(2) <sup>f</sup>	5.6807(3)	11.5243(6)	8.2515(6)	540.20(5)	3338	Holl et al. [2008]
Fe00	0	0 <sup>d</sup>		5.6921(2)	11.460(1)	8.253(2)	538.3(2)	3472	Finger et al. [1993]
Fe08	0.08	0 <sup>d</sup>		5.7037(9)	11.4681(8)	8.2679(9)	540.1(1)	3585	Finger et al. [1993]
Fe16	0.16	0 <sup>d</sup>		5.7119(9)	11.4681(8)	8.2799(9)	542.4(1)	3693	Finger et al. [1993]
Fe25	0.25	0 <sup>d</sup>		5.717(1)	11.506(1)	8.299(1)	545.9(2)	3807	Finger et al. [1993]
Fe40	0.40	0 <sup>d</sup>		5.739(2)	11.515(2)	8.316(1)	549.6(2)	4010	Finger et al. [1993]
Z0902	0.10	0.25 <sup>e</sup>		5.7045(4)	11.4765(4)	8.2707(4)	541.46(5)	3595	This study
Z0570	0.11	1.4 <sup>e</sup>	2.0(2) <sup>f</sup>	5.6929(2)	11.5275(2)	8.2648(1)	542.37(2)	3515	This study

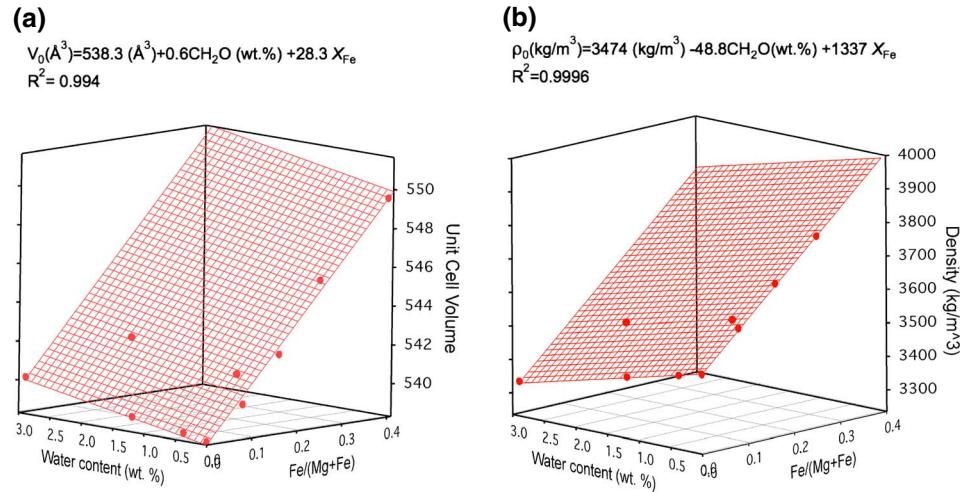
<sup>a</sup>Calibration of Paterson [1982].<sup>b</sup>Calibration of Libowitzky and Rossman [1997].<sup>c</sup>X-ray diffraction,  $b/a$  axial ratio method of Jacobsen et al. [2005].<sup>d</sup>Not measured, assumed nominally anhydrous.<sup>e</sup>Current study, calibration of Deon et al. [2010].<sup>f</sup>Current study, NanoSIMS measurements.

The discrepancy between the H<sub>2</sub>O content determined by various FTIR calibrations and by NanoSIMS appears negligible when the water content is lower than ~0.5 wt % (Figure 6). However, for higher H<sub>2</sub>O content we observe a factor of about 1.5 difference between SIMS and FTIR, with SIMS giving systematically higher water contents, even when compared with the calibration of Deon et al. [2010] for sample Z0570. One reason why SIMS could show higher H<sub>2</sub>O contents would be due to small inclusions of very hydrous phases, such as superhydrous phase B. However, FTIR measurements would reveal such inclusions, which have not been observed in any of the wadsleyite crystals. It is more likely that FTIR spectra from the previous studies underestimated the H<sub>2</sub>O content, either because polarized measurements did not use a mineral-specific calibration, because unpolarized measurements did not sample enough directions to build a statistically significant sampling of total absorbance, or because in very hydrous samples the main O-H stretching band of wadsleyite completely absorbed the light.



**Figure 6.** H<sub>2</sub>O content of five wadsleyite samples measured by SIMS (this study), in comparison with water contents reported by FTIR in original studies: WS3056 from Jacobsen et al. [2005] and Holl et al. [2008] is nearly anhydrous with only 0.005–0.006 wt % H<sub>2</sub>O as measured by both FTIR and SIMS (Table 4); WH833 from Jacobsen et al. [2005] and Mao et al. [2008] used the polarized-FTIR calibration of Libowitzky and Rossman [1997]; WH2120 and SS0401 from the study of Mao et al. [2008] from  $b/a$  axial ratio method [Jacobsen et al., 2005], which is tied largely to FTIR calibrations of Paterson, [1982] and Libowitzky and Rossman [1997]; Z0570 from the current study gives 1.2 wt % H<sub>2</sub>O using Libowitzky and Rossman [1997] or 1.4 wt % H<sub>2</sub>O (as plotted) using the FTIR calibration of Deon et al. [2010]. The SIMS measurement of Z0570 was conducted on six different crystals, up to 14 points per crystal, using two different ion probes and calibration methods (see Table 2 for details).

Regression analyses of the effect of Fe and H<sub>2</sub>O on the lattice parameters of wadsleyite are shown in Figure 5. Incorporation of 1 wt % H<sub>2</sub>O in the Mg end-member wadsleyite causes the  $a$  and  $c$  axes of the wadsleyite to contract by ~0.001 and 0.005 Å, respectively, whereas the  $b$  axis expands by 0.036 Å. As seen in Table 4, uncertainty in the individual lattice parameters is typically in the fourth decimal place. The influence of 0.10  $X_{\text{Fe}}$  on  $b_0$  is only half that of incorporating 1 wt % H<sub>2</sub>O into Mg-wadsleyite. In contrast, the influence of 0.10  $X_{\text{Fe}}$  on  $a_0$  and  $c_0$  is more significant than adding 1 wt % H<sub>2</sub>O (Figure 5).



**Figure 7.** (a) The unit cell volume of Mg- and Fe-wadsleyite samples listed in Table 4 as a function of both Fe and H<sub>2</sub>O content; (b) the variation of density ( $\rho_0$ ) with Fe and H<sub>2</sub>O content (in wt %). The blue meshes represent the planar regression fitting results.

To quantify the combined effects of Fe and H<sub>2</sub>O, we carried out a 3-D planar multivariable linear regression analysis using all of the data listed in Table 4. The substitution of Mg by Fe and incorporation of H<sub>2</sub>O into wadsleyite both expand the unit cell volume, as illustrated in Figure 7a. A 3-D fit to all of the data in Table 4 yields the following relationship for  $V_0$  of wadsleyite:

$$V_0(\text{Å}^3) = 538.3(\pm 0.6)(\text{Å}^3) + 0.6(\pm 0.2)C_{\text{H}_2\text{O}}(\text{wt \%}) + 28.3(\pm 1.7)X_{\text{Fe}}, \quad (1)$$

with a regression fitting parameter of  $R^2 = 0.9936$  and  $\chi^2 = 0.848$ . The variation of initial density ( $\rho_0$ ) with Fe and H<sub>2</sub>O content (in wt %) is illustrated in Figure 7b with fitted relationship:

$$\rho_0(\text{kg/m}^3) = 3474(\pm 4)(\text{kg/m}^3) - 49(\pm 3)C_{\text{H}_2\text{O}}(\text{wt \%}) + 1337(\pm 24)X_{\text{Fe}}, \quad (2)$$

with regression fitting parameter of  $R^2 = 0.9996$  and  $\chi^2 = 169.154$ . Variation of the reference density of wadsleyite with both Fe and H<sub>2</sub>O content will serve as an important anchor for high-pressure equations of state in this system.

### 3.2. Influence of Fe and H<sub>2</sub>O on the P-V Equation of State of Wadsleyite

To investigate combined influence of Fe and H<sub>2</sub>O on the P-V equation of state of wadsleyite, unit cell volumes of sample Z0902 and sample Z0570 at high pressures were collected and analyzed as described in section 2. The resulting pressure-volume data of wadsleyite Z0902, wadsleyite Z0570, ringwoodite Z0570, and MgO are listed in Table 5. As shown in Figure 8, the initial unit cell volume ( $V_0$ ) of sample Z0570 (2.0 wt % H<sub>2</sub>O) is about 0.14% larger than  $V_0$  of Z0902 (with 0.25 wt % H<sub>2</sub>O), but because the hydrous sample is more compressible, a volume crossover occurs at about 6 GPa. Although the hydrous sample has a smaller unit cell volume above 6 GPa, it remains less dense across the experimental pressure range (Figure 9) because of its lower atomic mass (Table 3). The density difference between hydrous (Z0570) and the nominally anhydrous (Z0902) wadsleyite is ~2.8%, a relationship which remains unchanged up to the maximum pressure of the experiment at 32 GPa, suggesting a very similar value of  $K_0'$ .

Pressure-volume data were fitted to a second- and third-order Birch-Murnaghan equation of state (BM-EOS) using EoSFit7 [Angel *et al.*, 2014]. In our procedure,  $V_0$  and  $K_{T0}$  were refined for both second- and third-order BM equations. The value of  $K_0'$  was also refined for the third-order BM-EOS fitting, while  $K_0'$  has an implied value of 4 in the second-order BM-EOS fitting. The uncertainties in pressure and volume measurements were used as weights. Prior to fitting, we tested the reliability of ruby pressures in this experiment by fitting the P-V data of MgO (Table 5) to a third-order BM-EOS. The resulting equation of state is  $V_0 = 74.69(2)\text{ Å}^3$ ,  $K_{T0} = 163(2)$  GPa with  $K_{T0}' = 4.05(21)$ , in excellent agreement with previous static compression, ultrasonic, Brillouin, and theoretical studies [e.g., Dorogokupets, 2010; Jaocbsen *et al.*, 2002; Jacobsen *et al.*, 2008; Kennett and Jackson, 2009; Mukherjee *et al.*, 2013; Speziale *et al.*, 2001; Tange *et al.*, 2009; Wu *et al.*, 2008; Zha *et al.*, 2000]. This result gives us confidence to use ruby pressures as the pressure marker in this study, although

**Table 5.** Volume-Pressure Data From Single-Crystal X-ray Diffraction of Samples Compressed in the Same Diamond Anvil Cell<sup>a</sup>

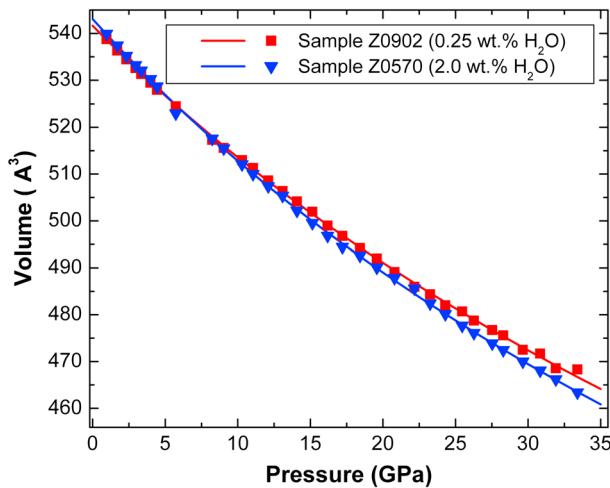
$P_{\text{ruby}}$ (GPa)	Wadsleyite (Z0902)	Wadsleyite (Z0570)	Ringwoodite (Z0570)	MgO
1.0	538.76(3)	539.92(2)	528.73(5)	74.33(1)
1.7	536.32(3)	537.41(2)	526.14(6)	73.97(1)
2.3	534.43(3)	535.19(2)	524.75(7)	73.58(1)
3.0	532.57(2)	533.21(2)	522.65(5)	73.40(1)
3.4	531.32(3)	532.05(1)	522.20(5)	73.23(1)
4.0	529.43(3)	530.22(1)	519.54(6)	72.91(1)
4.4	527.99(3)	528.65(1)	519.24(5)	72.74(1)
5.7	524.47(2)	--	515.88(6)	72.20(1)
6.3	--	522.97(2)	514.13(5)	72.05(1)
7.1	--	--	512.83(6)	71.71(1)
8.2	517.30(3)	517.59(2)	509.14(6)	71.25(1)
9.0	515.58(2)	515.48(2)	506.58(5)	70.98(1)
10.3	512.97(2)	512.09(2)	504.46(6)	70.57(1)
11.1	511.29(2)	510.01(2)	502.48(5)	70.39(1)
12.1	508.62(2)	507.41(2)	500.51(6)	69.99(1)
13.1	506.37(2)	505.36(2)	497.83(5)	69.67(1)
14.1	504.15(2)	502.18(2)	496.10(7)	69.27(1)
15.1	501.95(2)	499.55(2)	493.91(5)	68.93(1)
16.2	498.97(3)	496.80(2)	491.57(5)	68.72(1)
17.2	496.78(3)	495.45(2)	487.44(6)	68.27(1)
18.4	494.20(2)	492.60(2)	486.57(5)	67.93(1)
19.6	492.01(2)	490.09(2)	484.04(5)	67.79(1)
20.8	489.11(2)	487.85(2)	482.24(5)	67.56(1)
22.2	485.85(2)	485.61(2)	478.79(5)	66.94(1)
23.3	484.35(2)	482.40(2)	477.44(6)	66.55(1)
24.3	482.00(2)	480.13(2)	475.67(5)	66.44(1)
25.5	480.69(2)	477.66(2)	473.08(4)	66.15(1)
26.3	478.72(2)	476.10(2)	471.46(4)	65.81(1)
27.5	476.71(2)	473.87(2)	469.30(4)	65.76(1)
28.3	475.57(2)	472.46(2)	468.42(4)	--
29.6	472.50(2)	469.99(2)	465.75(4)	65.04(1)
30.8	471.69(2)	468.06(2)	464.21(4)	64.78(1)
31.9	468.56(2)	466.19(2)	462.34(4)	64.51(1)

<sup>a</sup>Unit cell volumes are in  $\text{\AA}^3$ . Uncertainty in the ruby pressure is  $\pm 0.1$  GPa based on  $P$  measured before and after measurement of sample volume.

we stress, because the volume of each sample was measured at the same pressure in a single experiment, we ultimately aim to determine if there is any difference in the fitted  $K_0'$  of hydrous versus anhydrous wadsleyite, independent of the secondary pressure scale.

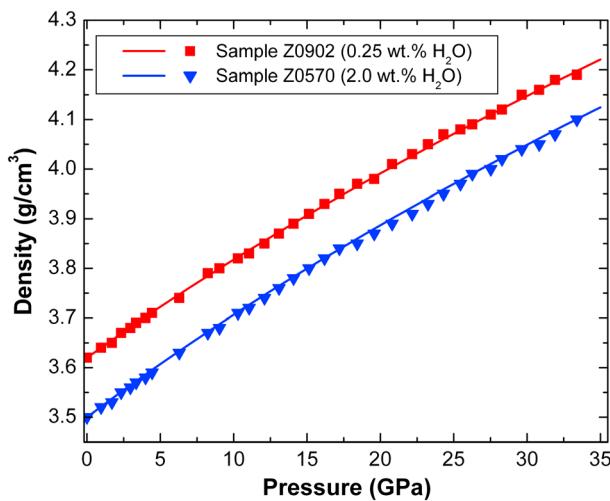
The resulting second-order BM-EOS parameters are  $V_0 = 542.0(2) \text{\AA}^3$ ,  $K_{T0} = 166(1) \text{ GPa}$  and  $V_0 = 543.4(2) \text{\AA}^3$ ,  $K_{T0} = 153(1) \text{ GPa}$  for slightly hydrous (Z0902) and very hydrous (Z0570) wadsleyite, respectively. The same pressure-volume data were fitted to the third-order BM-EOS, resulting in  $V_0 = 541.6(2) \text{\AA}^3$ ,  $K_{T0} = 170(2) \text{ GPa}$ , and  $K_{T0}' = 3.72(16)$  for slightly hydrous (Z0902) and  $V_0 = 543.1(3) \text{\AA}^3$ ,  $K_{T0} = 157(2) \text{ GPa}$ , and  $K_{T0}' = 3.77(14)$  for the very hydrous (Z0570) wadsleyite. The second-order BM-EOS obtained for hydrous ringwoodite (run Z0570) gives  $V_0 = 532.0(2) \text{\AA}^3$  and  $K_{T0} = 172(1)$ , with parameters  $V_0 = 531.6(3) \text{\AA}^3$ ,  $K_{T0} = 177(3) \text{ GPa}$ , and  $K_{T0}' = 3.7(2)$  for the third-order BM-EOS. All values of  $K_{T0}'$  resulting from this experiment are lower than the reported values of  $K_0'$  from previous studies, which range from 4.0 to 7.1 (Figure 1).

Two approaches were taken to test our fitting results. First, the same data sets were fitted to second- and third-order BM equations of state using the MgO pressure marker, resulting in EOS parameters that are identical within error to the results using ruby pressures. In order to determine whether the improvement in fit to the data achieved by using a third-order BM-EOS (in which the value of  $K_{T0}'$  is a free parameter) rather than second-order BM (in which  $K_{T0}'$  is by definition fixed at a value of 4) is statistically meaningful, we performed  $F$  test analyses for the addition of the free parameter [Bevington, 1969]. The results show that the deviation of  $K_{T0}'$  from the reference value of 4 is statistically significant at the 96% confidence level for sample Z0902 and



**Figure 8.** Volume-pressure data for anhydrous (red, Z0902) and hydrous (blue, Z0570) Fe-bearing wadsleyite and fitted equations of state (solid curves) given in the text. Although hydrous wadsleyite has a larger cell volume at room pressure, a crossover occurs at about 6 GPa due to the higher compressibility of the hydrous sample.

(shown in the inset in Figure 10) indicates that the  $V_0$  measured by the four-circle X-ray diffractometer is not compatible with the high-pressure volumes collected using the synchrotron X-ray diffractometer. We therefore excluded the measured four-circle  $V_0$  from our  $P$ - $V$  data and carried out the final least squares fitting following the method described in previous paragraphs. Because  $V_0$  is a prerequisite for calculating the Eulerian finite strain, the measured  $V_0$  was replaced with  $V_0$  obtained from the third-order BM equation of state to calculate the experimental Eulerian finite strain in the  $F_E$ - $f_E$  plot reported here. As seen in Figure 10, the EOS of both Z0902 and Z0570 are well constrained by the data and have slightly negative slopes in this  $F_E$ - $f_E$  plot, indicating that the values of  $K_{T0}'$  are smaller than 4.0. The regression lines for Z0902 and Z0570 data points are almost parallel to each other. Our results suggest that increasing  $H_2O$  from 0.25 to 2.0 wt %  $H_2O$  in Fe-bearing wadsleyite reduces the isothermal bulk modulus ( $K_{T0}$ ) from 170(2) GPa to 157(2) GPa, a reduction of 10.5%, but hydration does not affect the pressure derivative.

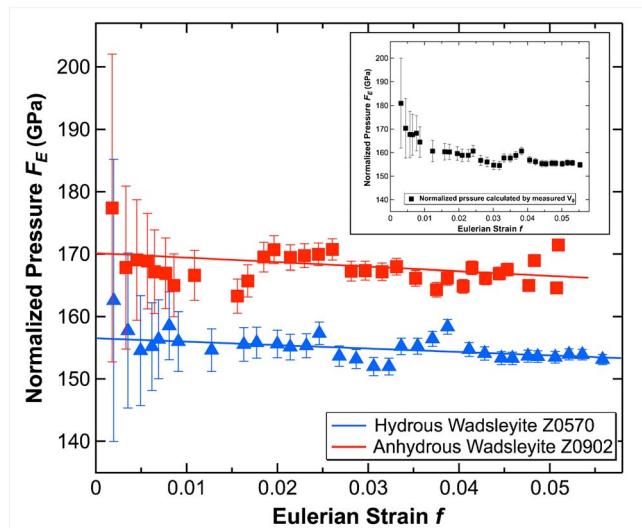


**Figure 9.** Variation of density with pressure for anhydrous (red, Z0902) and hydrous (blue, Z0570) Fe-bearing wadsleyite from the current compression experiments.

the 94% confidence level for both beta and gamma phases of Z0570. The low  $K_{T0}'$  values of Fe-bearing wadsleyite observed in this study are likely caused by the combined influences of Fe and  $H_2O$  in wadsleyite, rather than by experimental issues, as the resulting  $K_{T0}'$  of MgO within the same pressure chamber is 4.05(21), which is consistent with previous studies, e.g., 4.08(4) from the simultaneous nonlinear inversion of multiple data sets performed by Kennett and Jackson [2009].

To better visualize the quality of EOS fits, a plot of normalized stress  $F_E = P/3f_E(1+2f_E)^{5/2}$  versus Eulerian strain  $f_E = [(V_0/V)^{2/3} - 1]/2$  (an  $F_E$ - $f_E$  plot) was constructed as shown in Figure 10. The unusual curvature of the  $F_E$ - $f_E$  plot at very low strain for wadsleyite Z0570 calculated using the measured  $V_0$

Before discussing the combined influence of Fe and  $H_2O$  on the compressibility of wadsleyite, we briefly review the independent effects of Fe and  $H_2O$  on the compressibility of wadsleyite separately. Many studies have examined the effect of Fe on elastic properties of wadsleyite [e.g., Hazen et al., 2000a; Hazen et al., 2000b; Li and Gwanmesia, 1996; Li and Liebermann, 2000; Liu et al., 2009; Wang et al., 2014; Zha et al., 1997]. Experimental results suggest that Fe has an insignificant effect on the bulk modulus ( $K_{T,S0}$ ) of anhydrous wadsleyite: the bulk modulus of 0.0  $X_{Fe}$  wadsleyite is ~170(2) GPa [Li and Gwanmesia, 1996; Zha et al., 1997; Hazen et al., 2000a], 0.075  $X_{Fe}$  wadsleyite is 170(3) GPa [Wang et al., 2014], 0.12  $X_{Fe}$  wadsleyite is ~172(2) GPa [Li and Liebermann, 2000],

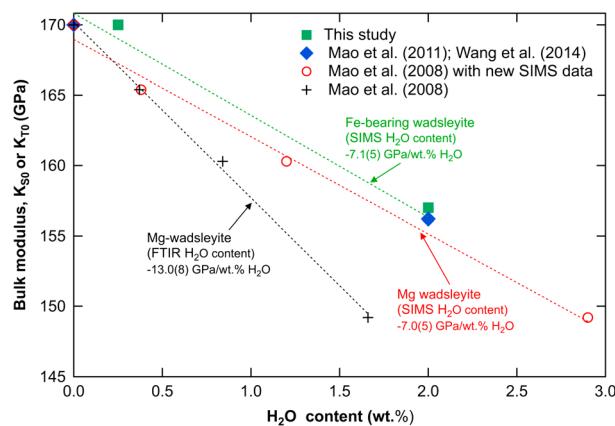


**Figure 10.** Plot of the normalized (Birch) pressure ( $F_E$ ) versus Eulerian strain ( $f$ ) for anhydrous (red) and hydrous (blue) wadsleyite. Data were calculated using the fitted  $V_0$  obtained from the third-order BM equations of state. The inset shows the  $F_E$  of sample Z0570 calculated using the measured  $V_0$  at ambient conditions. The curvature toward high  $F_E$  at very low strain indicates the measured  $V_0$  from a separate experiment is not compatible with the high-pressure data and therefore not used in the final equation of state fitting.

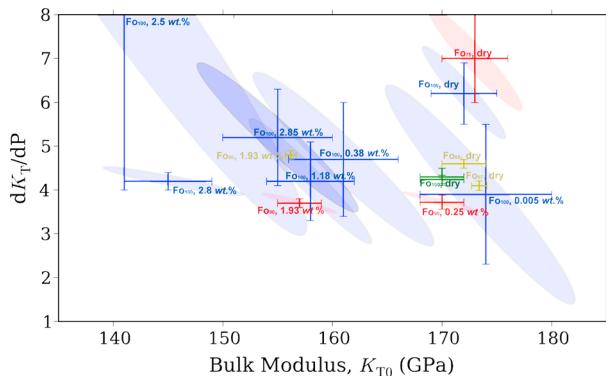
$H_2O$  causes a 7.6% decrease in  $K_{S0}$ . The fit to  $K_{S0}$  by Mao et al. [2008] was given as  $K_{S0} = 170.9(9) - 13.0(8)C_{H2O}$ , where  $C_{H2O}$  is the water content in weight percent. However, we note that the samples measured by Mao et al. [2008], namely, WH833, WH2120, and SS0401, are from the same synthesis runs that we have reanalyzed by NanoSIMS in the current study, showing considerably more  $H_2O$  in WH2120 and SS0401 than estimated by FTIR methods.

The effect of  $H_2O$  on the bulk modulus of Mg- and Fe-bearing wadsleyite is illustrated in Figure 11. We include data from both static compression (current study) and several previous Brillouin spectroscopy studies for comparison [Mao et al., 2008, 2011; Wang et al., 2014]. The relation between the adiabatic ( $K_{S0}$ ) and isothermal ( $K_{T0}$ ) bulk modulus is  $K_T = K_S/(1 + \alpha\gamma T)$ ; however, because the  $1 + \alpha\gamma T$  term is close to unity, for the purpose of the current discussion we will refer to  $K_{S0}$  and  $K_{T0}$  interchangeably as  $K_0$ .

In Figure 11, the variation of  $K$  with water content is most steep for the Mg-wadsleyite data of Mao et al. [2008], with  $dK_0/dC_{H2O} = -13.0(8) \text{ GPa}/\text{wt \% } H_2O$ . However, when we apply the new NanoSIMS  $H_2O$  contents measured in this study on samples from the same synthesis runs, the slope decreases dramatically to  $dK_0/dC_{H2O} = -7.0(5) \text{ GPa}/\text{wt \% } H_2O$  (Figure 11). For Fe-bearing wadsleyite, including the anhydrous value from Wang et al. [2014], the two Fe-bearing samples from this study, and the Brillouin measurements of SZ0570 from Mao et al. [2011], we



**Figure 11.** Bulk modulus ( $K_0$ ) versus  $H_2O$  content for wadsleyite. The black dashed line represents the fitting result with FTIR  $H_2O$  content after Mao et al. [2008], and the red dashed line represents the new result with SIMS  $H_2O$  values. The bulk modulus versus  $H_2O$  relationship for Fe-bearing wadsleyite is presented by green dashed line fit to Wang et al. [2014], Mao et al. [2011], and the current study.



**Figure 12.** Confidence ellipsoids constructed from data sets obtained in this and previous studies. References used in this figure are listed in Table 6. The blue ellipses represent Mg end-member wadsleyite with different  $H_2O$  content. The red ellipses represent Fe-bearing wadsleyite with different  $H_2O$  content. The equations of state measured by sound velocity techniques are shown in deep and light green symbols.

oxygen atoms in the M3 vacancies. However, the current study using the comparative compressibility method on Fe-bearing wadsleyite does not show an effect of  $H_2O$  on  $K_0'$ . The  $K_{T0}'$  of Z0902 (0.25 wt %  $H_2O$ ) is indistinguishable within mutual uncertainties from the  $K_{T0}'$  of Z0570 (2.0 wt %  $H_2O$ ) in this study. Fitting the  $P$ - $V$  data in XRD compression experiments suffers from a well-known correlation issue between fitted  $K$  and  $K_0'$ . In order to illustrate this correlation, we refitted the EOS of previously reported data and constructed confidence ellipses for each data set as shown in Figure 12. Reconstructing confidence ellipses of previous data allows us to compare not only the best fitting third-order BM equations of state but also the potential combinations of ( $K_{T0}$ ,  $K_{T0}'$ ) which should fit the data equally well within the 63% confidence region ( $1\sigma$ ) as the best fitting results. For comparison (Figure 12), results from sound velocity measurements were also converted to  $K_{T0}$  and  $K_{T0}'$  via thermodynamic relationships. Figure 12 demonstrates that the EOS is less well constrained in experiments with smaller pressure ranges, especially the value of  $K_{T0}'$  which is the second-order derivative of the  $P$ - $V$  curve and strongly affected by the pressure range of data. The  $K_{T0}'$  for hydrous Mg-wadsleyite is best constrained in the experiment performed by Ye *et al.* [2010]. The observed range of  $K_{T0}'$  for a very hydrous Mg-wadsleyite (2.8 wt %  $H_2O$ ) is between 4 and 5. Although Holl *et al.* [2008] suggested that increasing  $H_2O$  content elevates the value of  $K_{T0}'$ , this is not clearly seen from this plot. The confidence ellipse of Z0570 is almost identical to the confidence ellipse of Z0902, only shifted to lower  $K_{T0}$ . This confirms that the combined influence of Fe and  $H_2O$  does not affect  $K_{T0}'$ . Therefore, more hydrated Fe-bearing wadsleyite becomes more compressible and remains more compressible than anhydrous wadsleyite at transition zone pressures.

### 3.3. Comparison of Static Compression and Sound Velocity Measurements

As shown in Figure 12, the bulk modulus of wadsleyite samples with similar compositions is consistent across various experimental studies. Despite the remarkable variation in  $K_0'$  between studies (Figure 1), values from acoustic measurements are regarded to be more reliable than  $K_0'$  fitted from  $P$ - $V$  data because of the well-known correlation between  $K$  and  $K_0'$  (illustrated in Figure 12). The elastic properties of the hydrous Fe-bearing wadsleyite sample used in this study (run Z0570) were also studied by Mao *et al.* [2011] using Brillouin spectroscopy up to 12 GPa, allowing a direct comparison. The bulk modulus obtained from this study and Mao *et al.* [2011] is consistent within mutual uncertainties (Table 6). However, the  $K_0'$  from Mao *et al.* [2011] of 4.8(1) falls considerably outside the confidence ellipse of the sample Z0570 in this study with  $K_0' = 3.8(1)$  (Figure 12). This suggests that these two data sets are incompatible. In order to examine the source of the discrepancy in  $K_0'$  between our study and Mao *et al.* [2011], we compared the  $V/V_0$  ratios at high pressures in the two experiments. The  $V/V_0$  ratio in the sound velocity measurement was calculated from the  $\rho_0/\rho$  ratio, where density  $\rho_0$  was calculated using the measured  $V_0$ , and  $\rho$  values at high pressures were estimated via a “self-consistent method” [Mao *et al.*, 2011]. In this method, an initial density model was used to obtain  $K_{S0}$  and  $K_{S0}'$  from a least square fit to the finite strain

obtain  $dK_0/dC_{H2O} = -7.1(5)$  GPa/wt %  $H_2O$ . Thus, our new interpretation of all the data is that water influences the bulk modulus of Fe-free and Fe-bearing wadsleyite similarly, with an offset to slightly lower values for Fe-bearing samples due to Fe substitution for Mg.

In contrast to the consistent reduction in bulk modulus with increasing  $H_2O$  content across many experimental studies, the effect of  $H_2O$  on  $K_0'$  is not as systematic or straightforward (Figure 1). As noted earlier, reports range from 4 to  $>5$  from different studies on similar compositions. Compressibility studies in particular have suggested elevated  $K_{T0}'$  for hydrous Mg-wadsleyite [Holl *et al.*, 2008]. This was attributed to the repulsive force between

**Table 6.** Compilation of Equation of State Data for Wadsleyite and Ringwoodite<sup>a</sup>

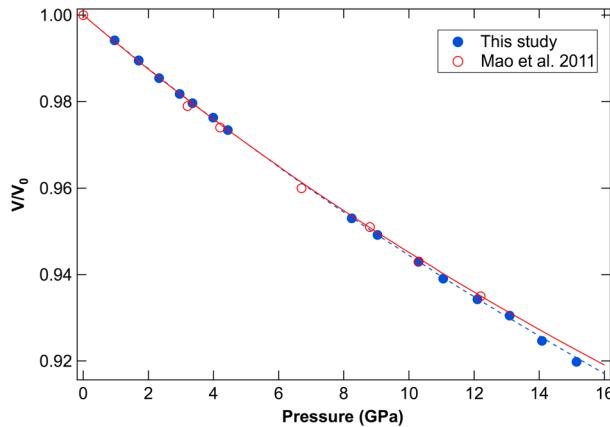
	$X_{\text{Fe}}$	$\text{H}_2\text{O}$ (wt %)	$K_0$ (GPa)	$K_0' = (\text{d}K/\text{d}P)_{P=0}$	Method <sup>b</sup>	$P_{\text{max}}$ (GPa), $T_{\text{max}}$ (K)	References
Anhy-wadsleyite	0	0	170(2)	4.24(10)	UI	12.5, RT	(1)
	0	0	169.2 (fixed)	4.1(1)	XRD	20, 2100	(2)
	0	0	170(2)	4.3(2)	BS	14, RT	(3)
	0	0	172(3)	6.3(7)	XRD	10.12, RT	(4)
	0.075	0	170(3)	4.1(1)	BS	17.7, RT	(17)
	0.12	0	172(3)	4.6(1)	UI	10, RT	(5)
	0.13	0	175.4(7)	4.10(11)	UI	12, 1073	(6)
	0.13	0	171.3(fixed)	4.74(8)	UI	12, 1073	(6)
	0.16	0	174(3)	4.0(fixed)	XRD	26, 900	(12)
	0.25	0	173(3)	7.1(8)	XRD	10.12, RT	(4)
Hyd-wadsleyite	0	0.005	173(5)	4.1(15)	XRD	7.3, RT	(7)
	0	0.37 (0.38 <sup>c</sup> )	161(4)	5.4(11)	XRD	9.01, RT	(7)
	0	0.37 (0.38 <sup>c</sup> )	165.4(9)	--	BS	Ambient	(8)
	0	0.84 (1.20 <sup>c</sup> )	160.3(7)	--	BS	Ambient	(8)
	0	1.18	158(4)	4.2(9)	XRD	8.56, RT	(7)
	0	1.66 (2.9 <sup>c</sup> )	154(4)	4.9(11)	XRD	9.58, RT	(7)
	0	1.66 (2.9 <sup>c</sup> )	149.2(6)	--	BS	Ambient	(8)
	0	2.5	155(2)	4.3(fixed)	XRD	8.5, RT	(9)
	0	2.8	137(5)	4.5(3)	XRD	60, RT	(10)
	0.10	0.25	170(2)	3.72(16)	XRD	31.9, RT	This study
Ringwoodite	0.11	1.93 (2.0 <sup>c</sup> )	156.2(5)	4.8(1)	BS	12, RT	(11)
	0.11	1.93 (2.0 <sup>c</sup> )	157(2)	3.77(14)	XRD	31.9, RT	This study
	0.09	0	188(3)	4.1(3)	BS	16, 923	(13)
	0.11	0.89	177(4)	5.3(4)	UI	9, RT	(14)
	0.11	0.79	175(3)	6.2(6)	XRD	45, RT	(15)
	0.14	1.1	175(1)	4.0(1)	BS	16, 673	(16)
	0.17	1.38	177(3)	3.7(2)	XRD	31.9, RT	This study

<sup>a</sup>References: (1) Li and Gwanmesia, [1996], (2) Katsura et al. [2009], (3) Zha et al. [1997], (4) Hazen et al. [2000a], (5) Li and Liebermann [2000], (6) Liu et al. [2009], (7) Holl et al. [2008], (8) Mao et al. [2008], (9) Yusa and Inoue [1997], (10) Ye et al. [2010], (11) Mao et al. [2011], (12) Fei et al. [1992], (13) Sinogeikin et al., [2003], (14) Jacobsen and Smyth [2006], (15) Manghnani et al. [2005], (16) Mao et al. [2012], and (17) Wang et al. [2014].

<sup>b</sup>Methods: UI, ultrasonic interferometry; BS, Brillouin scattering; XRD, X-ray diffraction.

<sup>c</sup> $\text{H}_2\text{O}$  content remeasured by SIMS (this study).

equation. The measured adiabatic bulk modulus is converted to isothermal bulk modulus ( $K_7$ ) at various pressures to construct the isothermal compression curve. The densities at high-pressure conditions are then determined using the isothermal compression curve and the initial density. The results of  $V/V_0$  versus  $P$  from both experiments are plotted together in Figure 13 and show a good agreement between data sets. We therefore conclude that the discrepancy in  $K_0'$  is likely to have been introduced by the numerical fitting processes.



**Figure 13.** Normalized volume ( $V/V_0$ , where  $V_0$  is the initial volume) for hydrous Fe-bearing wadsleyite (Z0570, filled symbols, this study) and calculated from the density ratio from Mao et al. [2011], red open circles.

Bass et al. [1981] found that when the value of  $K_{70}$  yielded by the compression data is compatible with the acoustically determined values within the mutual uncertainties, the accuracy of  $K_{70}'$  obtained in a  $P$ - $V$  data set can be improved by fixing  $K_{70}$  at the acoustically determined value in the least squares calculation. We therefore refit the compression data with  $K_{70}$  constrained at the acoustically determined value by Mao et al. [2011] and refined only  $V_0$  and  $K_{70}'$ . The resulting value of  $K_{70}'$  is 3.88(4), which is in excellent agreement with our original result  $K_0' = 3.77(14)$ . The resulting  $V_0$  is also consistent with the  $V_0$  in our original third-order BM equation of state. It can be seen in Figure 13 that

the two  $V/V_0$  versus  $P$  curves calculated from the EOS of the two experiments deviate from each other above 12 GPa, while the data points at  $\sim 12$  GPa from two experiments are still in good agreement with each other. As shown in Figure 12,  $K_{T0}'$  is the second-order derivative of the  $P$ - $V$  curve and especially sensitive to the pressure range over which data are obtained. We consider the potential influence of pressure range on  $K_{T0}'$  and refit the compression data only to 12 GPa; the resulting  $K_{T0}'$  is 4.16(14), still in reasonable agreement with the original result. The elevated  $K_{S0}'$  in the acoustic measurement is possibly affected by the slightly higher  $V/V_0$  ratio at  $\sim 12$  GPa, and the lack of constraints at higher pressures. Scanning through our fitting results with various pressure ranges and methods, a  $K_{T0}'$  value close to 4.0(0.2) is robust and may be a good assumption for a  $\text{Fo}_{90}$ -wadsleyite, in excellent agreement with the recent results for  $K_{S0}' = 4.1(1)$  from Wang et al. [2014]. The discrepancies in  $K_0'$  between the XRD and acoustic measurements are more likely rooted in the calibration and fitting methods across experiments than in the physical properties of samples. Regardless of the uncertain absolute value of  $K_0'$ , both Mao et al. [2011] and our study conclude that hydration has no effect on the  $K_0'$  of Fe-bearing wadsleyite. A similar absence of effect of hydration on  $K_0'$  of Fe-bearing olivine was also suggested by Chen et al. [2011].

### 3.4. Influence of $dK/dP$ on Bulk Sound Velocity in the Transition Zone

Comparison between sound velocities derived from mineralogical models and observed seismic velocity is potentially an important tool for constraining the hydration state of the transition zone. In order to demonstrate the importance and influence of  $K_0'$  on the predicted sound velocity of wadsleyite, the bulk sound velocity ( $V_\Phi = [K/\rho]^{1/2}$ ) of anhydrous and hydrous Fe-bearing ( $\text{Fo}_{90}$ ) wadsleyite was calculated using two self-consistent data sets (Table 7); one using  $K_0' = 3.8$  implied from the current study and one using  $K_0' = 4.8$  implied from the study of Mao et al. [2011]. The bulk sound velocity gradient for the  $K_0' = 3.8$  data set is 1.72 (m/s/km) for anhydrous wadsleyite and 1.81 m/s/km for hydrous wadsleyite in the depth range 410–520 km (Figure 14). Raising  $K_0'$  from 3.8 to 4.8 increases the velocity gradient by about 0.55 m/s/km, or about 30%, for both dry and hydrous wadsleyite. As shown in the current study, the effect of hydration on the velocity gradient in the wadsleyite-rich part of the transition zone is not expected to be significant due to the minimal effect of  $\text{H}_2\text{O}$  on  $K_0'$  (whether or not  $\sim 4$  or  $\sim 5$  is used). Instead, the major effect of hydration on Fe-bearing wadsleyite will be the reduction in sound velocities associated with the  $\sim 7$  GPa/wt%  $\text{H}_2\text{O}$  reduction in the bulk modulus, shown in Figure 11. Whether  $K_0' \sim 4$  or  $\sim 5$  is assumed, the difference in bulk sound velocity between anhydrous and hydrous wadsleyite is around 200 m/s, or 2.5%, at transition zone conditions (Figure 14). Accordingly, the expected reduction in bulk sound velocity in very hydrous regions of the transition zone would be around 1.5% in pyrolite (60% olivine component), which is potentially detectable seismically in high-resolution, regional studies.

In contrast to hydrous versus anhydrous olivine ( $\text{Fo}_{100}$ ) [Mao et al., 2010], comparing the bulk sound velocities of anhydrous and hydrous Fe-bearing wadsleyite ( $\text{Fo}_{90}$ ), no velocity crossover was observed in either the  $K_0' = 3.8$  or the  $K_0' = 4.8$  data sets. The differences in velocities between anhydrous and hydrous Fe-bearing wadsleyite are about 175 m/s (2.5%) and 165 m/s (2.2%) at 1673 K for  $K_0' = 3.8$  and  $K_0' = 4.8$ , respectively. These differences remain essentially constant in the depth range 410–520 km for both cases. However, if we remove the constraint imposed from this study that  $K_0'$  of anhydrous and hydrous wadsleyite are similar and instead compare  $V_\Phi$  for anhydrous wadsleyite with  $K_0' \sim 4$  [Katsura et al., 2009; Wang et al., 2014] and  $V_\Phi$  for hydrous wadsleyite with  $K_0' = \sim 5$  [Holl et al., 2008; Mao et al., 2011], a velocity crossover occurs at a depth of  $\sim 350$  km, implying that hydrous wadsleyite would be the faster phase at transition zone pressures, which seems highly unlikely.

The importance of constraining  $K_0'$  also becomes evident in attempting to estimate of the fraction of olivine in Earth's mantle. Assuming the Fe content in olivine and wadsleyite in the Earth's mantle is about 10 mol % and  $\text{H}_2\text{O}$  partitioning of 1:5 between olivine and wadsleyite [Inoue et al., 2010], the contrast in bulk sound velocity ( $\Delta V_{\alpha-\beta}$ ) at 410 km along a 1673 K adiabat calculated with thermal parameters in Table 7 and using  $K_0 = 123(3)$  GPa and  $K_0' = 4.6$  for 0.4 wt%  $\text{H}_2\text{O}$  Fe-bearing olivine ( $\text{Fo}_{90}$ ) [Chen et al., 2011], we obtain  $\Delta V_{\alpha-\beta} = 4.7\%$  against hydrous wadsleyite with  $K_0' = 3.8$  and  $\Delta V_{\alpha-\beta} = 8.1\%$  when  $K_0' = 4.8$  is used. As a consequence, the bulk sound velocity contrast of 3.3% used in the IASP91 model [Kennett and Engdahl, 1991] would correspond to around 70% olivine under hydrous conditions when  $K_0' = 3.8$  is used, but only 40% olivine if  $K_0' = 4.8$  is used. A similar wide range of possible results is obtained for dry conditions (dry moduli for olivine and wadsleyite), yielding  $\sim 60\%$  olivine when  $K_0' = 3.8$  is used for anhydrous wadsleyite and  $\sim 37\%$

**Table 7.** Thermoelastic Data Sets Used to Compare Bulk Sound Velocity in the Transition Zone Under Anhydrous (Anhy) and Hydrous (Hyd) Conditions for Various Scenarios of  $K_0'$  Discussed in the Text<sup>a</sup>

	$K_{S0}$ (GPa) <sup>b</sup>	$K_{S0}'$	$\rho_0$ (g/cm <sup>3</sup> )	$dK_S/dT$ (GPa K <sup>-1</sup> )	$\alpha \times 10^{-5}$ (K <sup>-1</sup> )	Reference
<i>Data Set 1</i>						
Anhy-wadsleyite	172(2)	3.72(16)	3.600	-0.0170	3.40	d
Hyd-wadsleyite	158(2)	3.77(14)	3.520	-0.0170	3.40	d
Anhy-olivine	129(4)	4.6(fixed)	3.351 <sup>c</sup>	-0.0175	3.64	e
Hyd-olivine	124(3)	4.6(fixed)	3.331 <sup>c</sup>	-0.0175	3.64	e
<i>Data Set 2</i>						
Anhy-wadsleyite	171.3(fixed)	4.74(8)	3.633	-0.0170	3.40	f
Hyd-wadsleyite	156.2(5)	4.8(1)	3.520	-0.0170	3.40	g
Anhy-olivine	131.1(1.9)	3.8(2)	3.343	-0.0175	3.64	h
Hyd-olivine	124(3)	3.8(2)	3.331	-0.0175	3.64	i

<sup>a</sup>For hydrous conditions, 0.4 wt % H<sub>2</sub>O in olivine and 2.0 wt % H<sub>2</sub>O in wadsleyite are used.

<sup>b</sup> $K_{T0}$  from this study and Chen et al. [2011] were converted to  $K_{S0}$  using  $K_{S0}/K_{T0} = 1.01$ .

<sup>c</sup>Calculated using the chemical compositions and  $V_0$  from Chen et al. [2011].

<sup>d</sup>This study, Liu et al. [2009], and Inoue et al. [2004].

<sup>e</sup>Chen et al. [2011], Liu and Li, [2006], Isaak et al. [2010], and Ye et al. [2009].

<sup>f</sup>Liu et al. [2009] and Inoue et al. [2004].

<sup>g</sup>Mao et al. [2011], Liu et al. [2009], and Inoue et al. [2004].

<sup>h</sup>Mao et al. [2011], Liu et al. [2009], and Inoue et al. [2004].

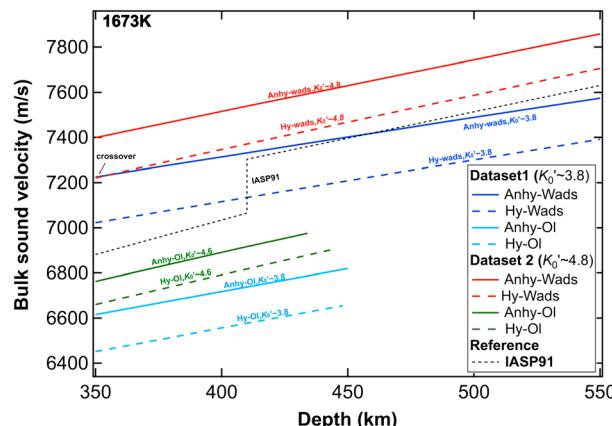
<sup>i</sup> $K_{S0}$  from Chen et al. [2011];  $K_{S0}'$  assumed equal to Anhy-olivine.

when  $K_0' = 4.8$  is used. Further, if  $K_0'$  of Fe-bearing olivine is reduced from 4.6 to 3.8, using a value of  $K_0 = 130$  GPa (Table 7) would give 38% and 28% olivine fraction when  $K_0' = 3.8$  and 4.8 for wadsleyite, respectively. However, the olivine fraction estimated using the self-consistent data sets (Table 7) is ~40% for both dry and hydrous conditions. Here we have demonstrated how it would be nearly impossible to accurately estimate the olivine fraction of the mantle near 410 km without better experimental constraints on  $K_0'$ .

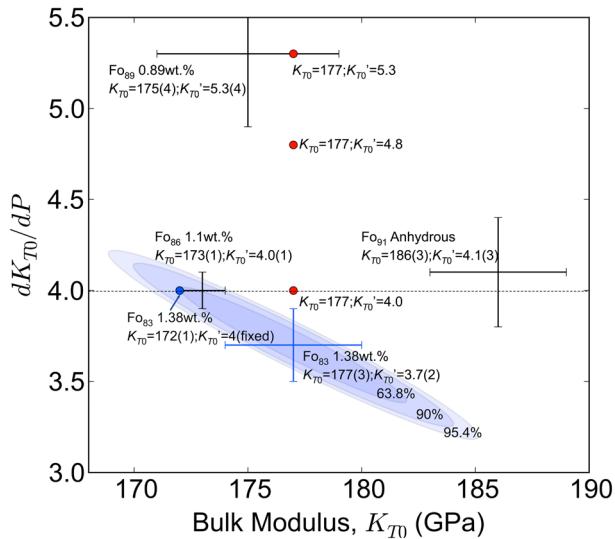
In this study, we also constrain the relative value of  $K_0'$  between Fe-bearing hydrous wadsleyite (Fo<sub>90</sub>) and ringwoodite (Fo<sub>83</sub>). The resulting  $K_0'$  of hydrous Fo<sub>83</sub> ringwoodite ( $X_{Fe} = 0.17$ , 1.38 wt % H<sub>2</sub>O) is 3.7(2), which is indistinguishable from the  $K_0' = 3.77(14)$  of the coexisting Fo<sub>90</sub> wadsleyite. The reported  $K_0'$  of hydrous Fe-bearing ringwoodite from previous studies ranges from 4.0(1) to 6.2(6) (Table 5 and Figure 15). We stress that in this study, we are more focused on accurately constraining the possible change in  $K_0'$  with composition (specifically, H<sub>2</sub>O content) and structure (between hydrous wadsleyite and hydrous ringwoodite), rather than on absolute values, which are subject to statistical trade-off with  $K_0$  in P-V fitting studies, as discussed above. Mao et al. [2012] suggested that hydration does not affect  $K_0'$  of Fe-bearing ringwoodite. The reported  $K_0' = 4.0(1)$  of a Fo<sub>89</sub> ringwoodite (1.1 wt % H<sub>2</sub>O) in Mao et al. [2012] is indistinguishable from  $K_0' = 4.1(3)$  reported for dry Fo<sub>91</sub> ringwoodite [Sinogeikin et al., 2003]. The values of  $K_0' \sim 4.0$  for dry and hydrous Fe-bearing ringwoodite from

Brillouin scattering [Mao et al., 2012; Sinogeikin et al., 2003] are significantly lower than the reported  $K_0' = 4.8(1)$  of hydrous Fo<sub>90</sub>-wadsleyite from Mao et al. [2011], implying a difference of 20% between  $K_0'$  of hydrous wadsleyite and hydrous ringwoodite. Ultrasonic measurements by Jacobsen and Smyth [2006] reported an extremely elevated  $K_{S0}' = 5.3(4)$  for hydrous Fo<sub>89</sub> ringwoodite (0.89 wt % H<sub>2</sub>O). We have found through comparative compressibility studies that the difference between  $K_0'$  of hydrous wadsleyite and hydrous ringwoodite is indistinguishable.

Next we compare the calculated bulk sound velocity jump between wadsleyite and ringwoodite ( $\Delta V_{\beta-\gamma}$ ) at transition

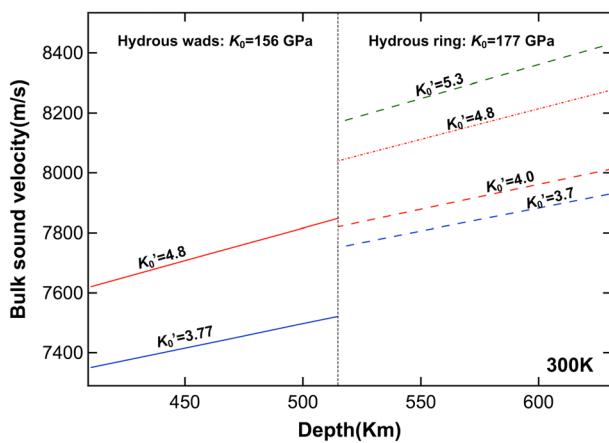


**Figure 14.** Comparison of bulk sound velocities for dry and hydrous phases along a 1673 K adiabat using  $K_0' = 3.8$  and  $K_0' = 4.8$ . Data sets 1 and 2 are defined in Table 7.



**Figure 15.** Confidence ellipsoids calculated from the equation of state fit for ringwoodite (Z0570) in this study. References for Fe-bearing ringwoodite used in this figure are listed in Table 6. The blue dot represents the result of a second-order BM equation of state for ringwoodite (Z0570). The equations of state measured by sound velocity techniques are converted to isothermal bulk moduli using the thermodynamic relationship in the text. The red dots represent the hypothetical equations of state used in Figure 15 to illustrate the effect of  $K_0'$  ranging from 4 to 5.3.

hydrous wadsleyite and hydrous ringwoodite to be the same as found in this study (either 3.8 or 4.8), the resulting  $\Delta V_{\beta-\gamma}$  is consistently around 2.7% (or 1.6% for 60% olivine component in pyrolite), consistent with the current experimental results. Whereas the absolute values of  $K_0'$  remain uncertain, the very tight constraint on the relative values of  $K_0'$  between phases from this study provide a very reasonable estimation for the predicted velocity jump at 520 km of around 1.5% under dry or hydrous conditions.



**Figure 16.** Comparison of bulk sound velocity jumps between hydrous wadsleyite and hydrous ringwoodite ( $\Delta V_{\beta-\gamma}$ ) calculated at 300 K using  $K_0 = 156$  GPa for wadsleyite and  $K_0 = 175$  GPa for ringwoodite, with various reported values of  $K_0'$  for ringwoodite. The  $\Delta V_{\beta-\gamma}$  calculated from this study is ~2.2%. The  $\Delta V_{\beta-\gamma}$  varies from -0.16% to 4.3% using values of  $K_0'$  that span the experimental literature. In the current study, we find no difference between  $K_0'$  of hydrous wadsleyite and hydrous ringwoodite, so for either  $K_0' \sim 4$  or  $\sim 5$ , the contrast is around 2.5% for wadsleyite, 1.5% for a transition zone composition of 60% olivine component.

zone pressures for cases in which  $K_0'$  of Fe-bearing ringwoodite is higher, equal, and lower than the  $K_0'$  of Fe-bearing wadsleyite. The equations of state for hydrous Fe-bearing wadsleyite (Z0570) from this study and Mao *et al.* [2011] were used in these calculations. For hydrous Fe-bearing ringwoodite, we used the  $K_0$  of ringwoodite (Z0570) from this study and varied  $K_0'$  from 4.0 to 5.3 to reflect the published range, as shown by the red dots in Figure 15. The bulk sound velocity jump ( $\Delta V_{\beta-\gamma}$ ) calculated at 300 K using the measured equations of state for coexisting wadsleyite and ringwoodite in this study is ~2.2% (Figure 16), corresponding to about 1.3% contrast for a 60%  $(\text{Mg},\text{Fe})_2\text{SiO}_4$  mantle component. If a value of  $K_0' = 4.8$  for hydrous wadsleyite is used [Mao *et al.*, 2011], the bulk sound velocity contrast with ringwoodite ( $\Delta V_{\beta-\gamma}$ ) ranges from -0.16% to 4.3% as  $K_0'$  of hydrous ringwoodite varies from 4.0 to 5.3 (Figure 16).

However, when we force the  $K_0'$  of

#### 4. Conclusions

Consistent trends for the effect of Fe and  $\text{H}_2\text{O}$  on  $K_0$  for wadsleyite and ringwoodite are in contrast to widely scattered values of  $K_0'$  found in recent experimental studies. Here we have discussed the importance of constraining  $K_0'$ , if only the difference in  $K_0'$  between compositions and structures of  $\text{Mg}_2\text{SiO}_4$ . Comparing the bulk sound velocity ( $V_\Phi$ ) of hydrous wadsleyite calculated with  $K_0' \sim 4$  versus  $K_0' \sim 5$  at a depth of 410 km results in a 230 m/s difference (3.3%), which is comparable to the velocity reduction caused by ~2 wt %  $\text{H}_2\text{O}$  in wadsleyite. Therefore, it is not feasible to accurately constrain the bulk mantle water content from current mineral physics equations of state parameters owing to the lack of consensus on  $K_0'$ . In this study, we provide strong evidence through comparative

compressibility studies that water does not influence  $K_0'$ , leaving the majority of influence expressed in reductions of  $K_0$ . We have revisited the composition of hydrous wadsleyite samples from a number of previous studies using SIMS and, in combination with the current results, revised estimates for the influence of  $H_2O$  on wadsleyite  $K_0$  from about  $-13$  GPa/wt %  $H_2O$  [e.g., Mao *et al.*, 2008] to  $-7$  GPa/wt %  $H_2O$  for both Mg- and Fe-bearing wadsleyite. The absolute value of  $K_0'$  from  $P$ - $V$  or acoustic methods is strongly affected by the numerical analysis and fitting processes. The resulting values of  $K_0'$  from recent studies are scattered even when the experimental data sets are consistent with each other, as we have demonstrated in this study. A more consistent data processing method across differing experimental techniques is needed for future studies.

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